

# A New Microscopic Multiphonon Approach to Nuclear Spectroscopy

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II*



# Semiclassical

$$\{\mathbf{a}_\lambda, \pi_\lambda\} \rightarrow \{\mathbf{O}_\lambda, \mathbf{O}_\lambda^\dagger\}$$

EoM

$$[H, \mathbf{O}_\lambda^\dagger] = \hbar\omega_\lambda \mathbf{O}_\lambda^\dagger$$

Collective  
modes

# Microscopic

TDA mapping

$$\mathbf{O}_\lambda^\dagger = \sum_{ph} \mathbf{c}_{ph}(\lambda) \mathbf{a}_p^\dagger \mathbf{a}_h$$

EoM

$$[H, \mathbf{O}_\lambda^\dagger] |> = \hbar\omega_\lambda \mathbf{O}_\lambda^\dagger |>$$

RPA mapping

$$\mathbf{O}_\lambda^\dagger = \sum_{ph} [X_{ph}(\lambda) \mathbf{a}_p^\dagger \mathbf{a}_h - Y_{ph}(\lambda) \mathbf{a}_h^\dagger \mathbf{a}_p]$$

EoM

$$[H, \mathbf{O}_\lambda^\dagger] |0> = \hbar\omega_\lambda \mathbf{O}_\lambda^\dagger |0>$$

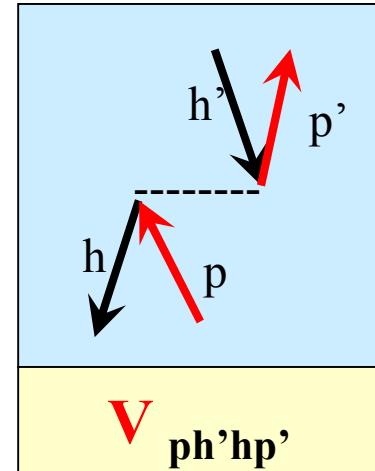
|0> ≡ correlated g.s

# TDA: Eigenvalue Equations

$$\langle \lambda | H | ph^{-1} \rangle = \langle \lambda | [H, a_p^\dagger a_h] | \rangle = \hbar \omega_\lambda c_{ph}^\lambda$$



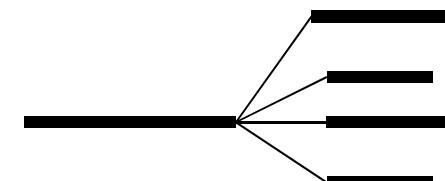
$$A c^\lambda = \hbar \omega_\lambda c^\lambda$$



$$A_{(ph)(p'h')} = (\epsilon_p - \epsilon_h) \delta_{pp'} \delta_{hh'} + V_{ph'h'p}$$

$$|\lambda\rangle = O_\lambda^\dagger | \rangle$$

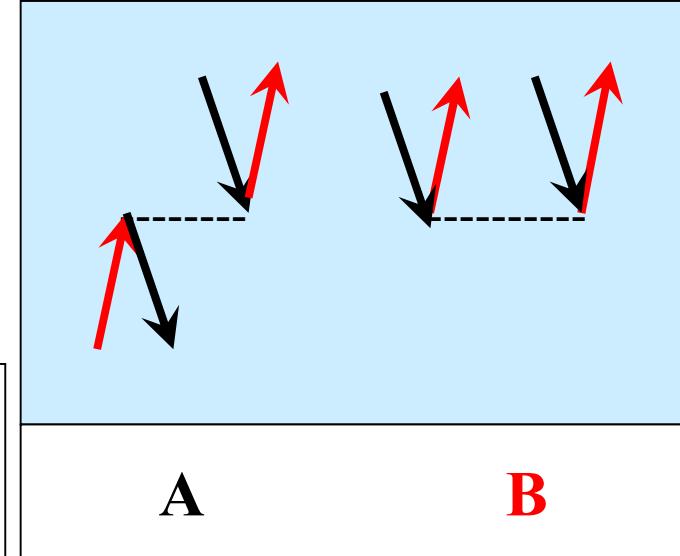
$$O_\lambda^\dagger = \sum_{ph} c_{ph}^\lambda a_p^\dagger a_h | \rangle$$



Landau damping

# RPA: Eigenvalue Equations

$$\begin{pmatrix} \mathbf{A} & \mathbf{B} \\ \mathbf{B} & \mathbf{A} \end{pmatrix} \begin{pmatrix} \mathbf{X}(\lambda) \\ \mathbf{Y}(\lambda) \end{pmatrix} = \hbar\omega_\lambda \begin{pmatrix} \mathbf{X}(\lambda) \\ -\mathbf{Y}(\lambda) \end{pmatrix}$$



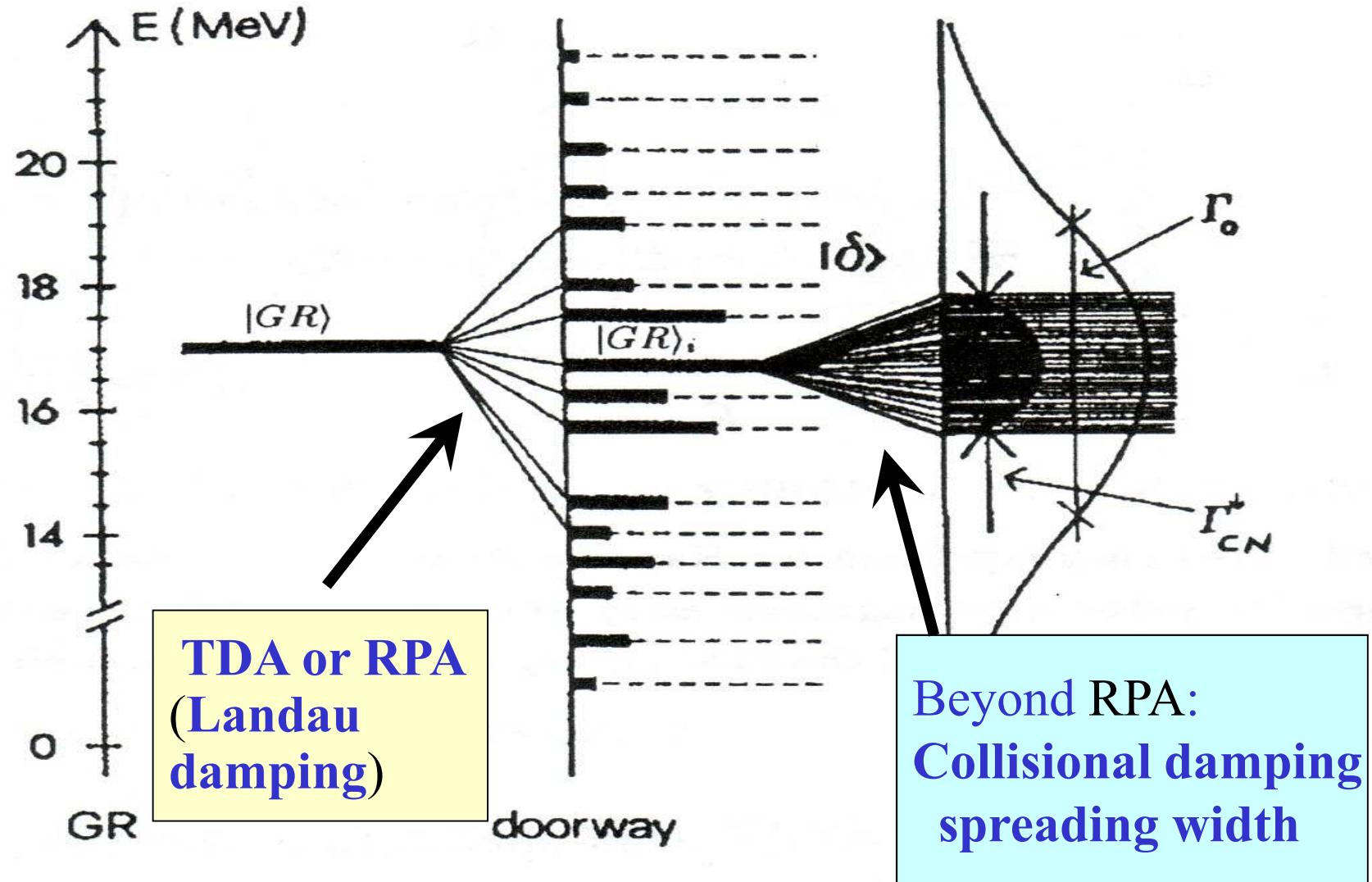
$$\begin{aligned} \mathbf{A}_{\text{ph}\mathbf{p}'\mathbf{h}} &= (\varepsilon_{\mathbf{p}} - \varepsilon_{\mathbf{h}}) \delta_{\mathbf{p}\mathbf{p}'} \delta_{\mathbf{h}\mathbf{h}'} + \mathbf{V}_{\text{ph}\mathbf{h}\mathbf{p}'}, \\ \mathbf{B}_{\mathbf{h}\mathbf{h}'\mathbf{p}\mathbf{p}'} &= \mathbf{V}_{\mathbf{h}\mathbf{h}'\mathbf{p}\mathbf{p}'} \end{aligned}$$

$$|\lambda\rangle = \mathbf{O}_\lambda^\dagger |0\rangle$$

$$\mathbf{O}_\lambda^\dagger = \sum_{\text{ph}} [\mathbf{X}_{\text{ph}}(\lambda) \mathbf{a}_\mathbf{p}^\dagger \mathbf{a}_\mathbf{h} - \mathbf{Y}_{\text{ph}}(\lambda) \mathbf{a}_\mathbf{h}^\dagger \mathbf{a}_\mathbf{p}] |0\rangle$$

Quasi-Boson Approximation:  $|0\rangle \rightarrow | \rangle$

# Collective modes: anharmonic features



From P.F. Bortignon, A. Bracco, R.A. Broglia,  
Giant Resonances (hap, 1998)

# Multiphonon excitations: Exp. evidence

## \* High-energy

(N. Frascaria, NP A482, 245c(1988);  
T. Auman, P.F. Bortignon, H.  
Hemling, Ann. Rev. Nucl. Part.  
Sc. 48, 351 (1998))

Double

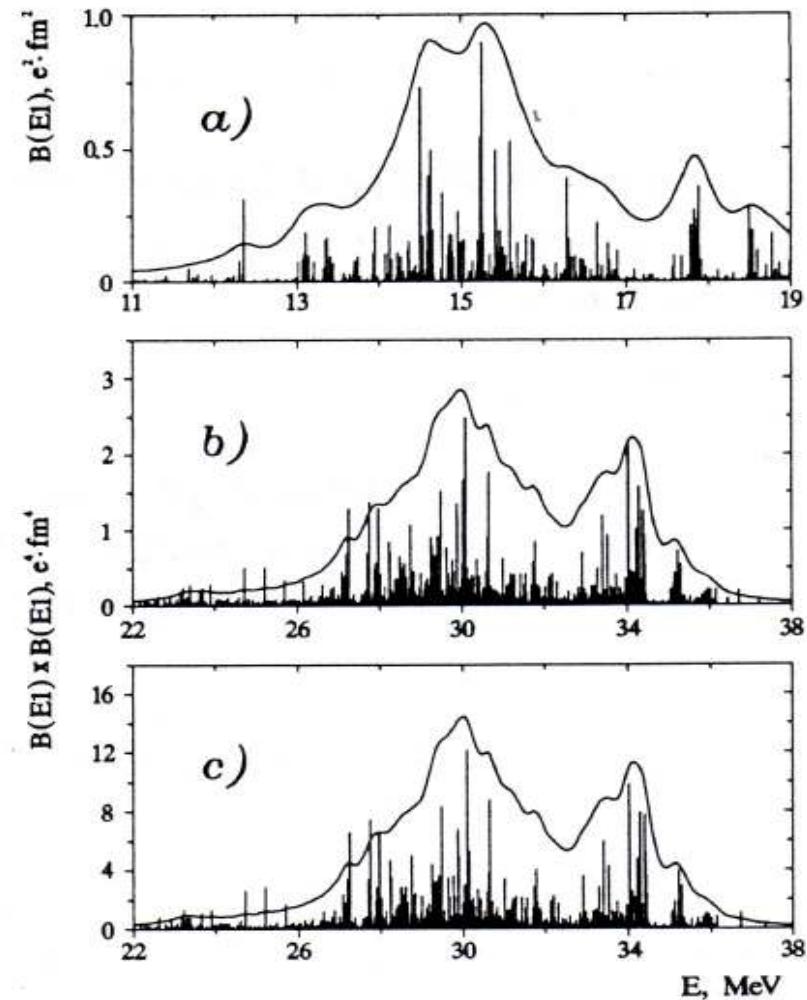
$$D \times D |0\rangle$$

and

triple

$$D \times D \times D |0\rangle$$

dipole giant resonances



# Multiphonon excitations: Exp. evidence

## \*\* Low-energy

M. Kneissl, H.H. Pitz, and A. Zilges, Prog. Part. Nucl. Phys. 37, 439 (1996); M. Kneissl, N. Pietralla, and A. Zilges, J.Phys. G, 32, R217 (2006) :

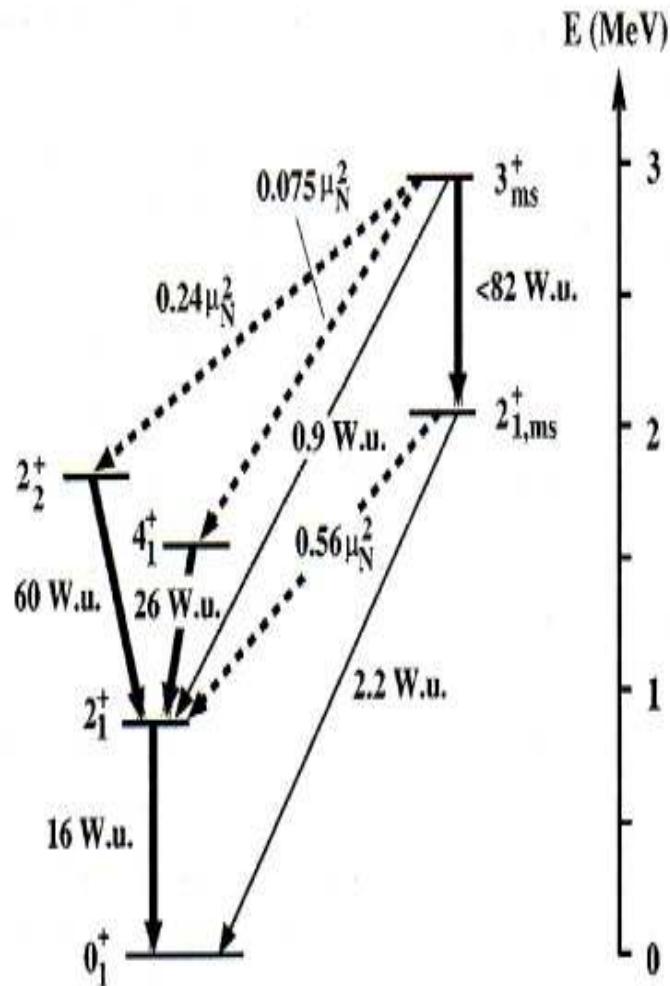
- Two- and three-phonon multiplets

$$Q_2 \times Q_3 |0\rangle, \quad Q_2 \times Q_2 \times Q_3 |0\rangle$$

- Proton-neutron (F-spin) mixed-symmetry states

(N. Pietralla et al. PRL 83, 1303 (1999))

$$[Q_2^{(p)} - Q_2^{(n)}] (Q_2^{(p)} + Q_2^{(n)})^N |0\rangle,$$



# From RPA to multiphonon approaches

Basic idea inspired by the **Boson expansion techniques (BET)**

(S. T. Belyaev and V. G. Zelevinsky, Nuc. Phys. 39, 582 (1962) )

$$O^\dagger \Rightarrow O^\dagger + O^\dagger O^\dagger + O^\dagger O^\dagger O^\dagger \dots$$

$$O^\dagger = \sum_{ij} [X(ij) \alpha_i^\dagger \alpha_j^\dagger - Y(ij) \alpha_i \alpha_j]$$

Recent calculations

a. Density functional approaches (up to two-phonons)

1. SRPA (Gambacurta et al. PRC8 1, 054312 (2010)) based on Skyrme

2. RRPA-phonon model (Litvinova et al. PRL 105, 022502 (2010))

b. A popular approach: QPM (Soloviev) (up to 3-phonons)

$$H(\text{separable}) \Rightarrow H(O^\dagger O) = \sum \omega_\lambda O_\lambda^\dagger O_\lambda + H_{vq}(O^\dagger O O^\dagger O)$$

Common features:

1. QBA is the underline approximation
2. No Correlations explicitly included in the gs

# Eigenvalue problem in multiphonon space

$$\mathbf{H} |\Psi_v\rangle = E_v |\Psi_v\rangle$$

$$|\Psi_v\rangle \in \mathcal{H} = \Sigma_n \oplus \mathcal{H}_n \quad \mathcal{H}_n \in |n; \beta\rangle \equiv n\text{-phonon basis states}$$

An obvious (**prohibitive!!**) choice

$$|n; \beta\rangle = O_{v_1}^\dagger \dots O_{v_i}^\dagger \dots O_{v_n}^\dagger |0\rangle$$

$$O_v^\dagger = \sum_{ph} c_{aph} a_p^\dagger a_h |0\rangle$$

A viable route

$$|n; \beta\rangle = \sum_{v\alpha} C_{v\alpha}^\beta O_v^\dagger |n-1; \alpha\rangle$$

# Construction of $|n; \beta\rangle$ : EofM

$$\langle n; \beta | [H, O_v^\dagger] | n-1; \alpha \rangle = (E_\beta^{(n)} - E_\alpha^{(n-1)}) \langle n; \beta | O_v^\dagger | n-1; \alpha \rangle$$



$$\sum_{v\gamma} [\mathcal{H}_{(\mu\alpha)(v\gamma)} - E_{\mathcal{D}} |_{(\mu\alpha)(v\gamma)}] C_{v\gamma} = 0$$

$$\mathcal{H} = \mathcal{AD}$$

$$\mathcal{D}_{(\mu\alpha)(v\gamma)} = \langle n, \gamma | O_v O_\mu^\dagger | n, \alpha \rangle$$

## 1-phonon: TDA

$$\mathbf{A} \mathbf{c} = \hbar\omega \mathbf{c}$$

$$\begin{aligned} \mathbf{A}_{(ph)(p'h')} &= (\varepsilon_p - \varepsilon_h) \delta_{pp'} \delta_{hh'}, \\ &+ \mathbf{V}_{ph'hp} \end{aligned}$$

$$|\lambda\rangle = \sum_{ph} \mathbf{c}_{ph}^\lambda \mathbf{a}_p^\dagger \mathbf{a}_h |>$$

## n-phonon: EofM

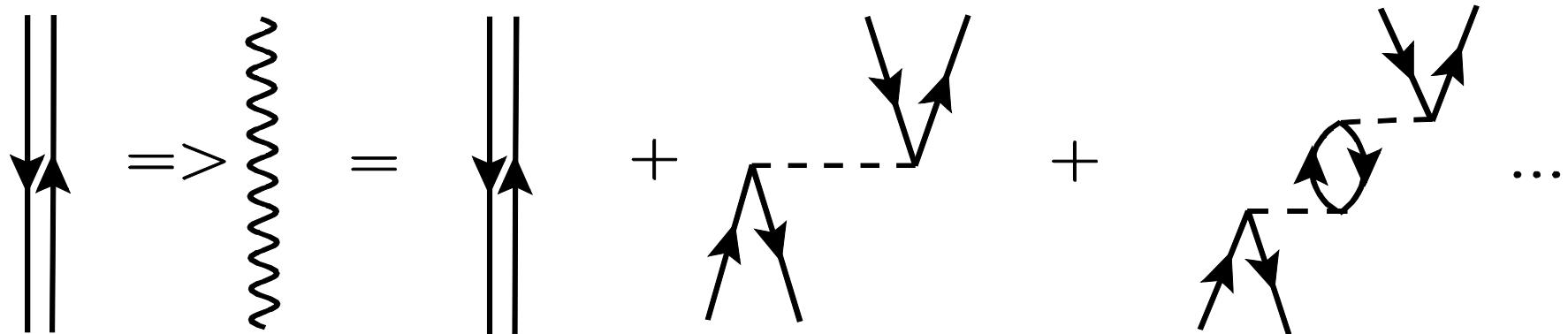
$$\mathcal{H}C = \mathcal{E} \mathcal{D} C$$

$$\mathcal{H} = \mathcal{A} \mathcal{D}$$

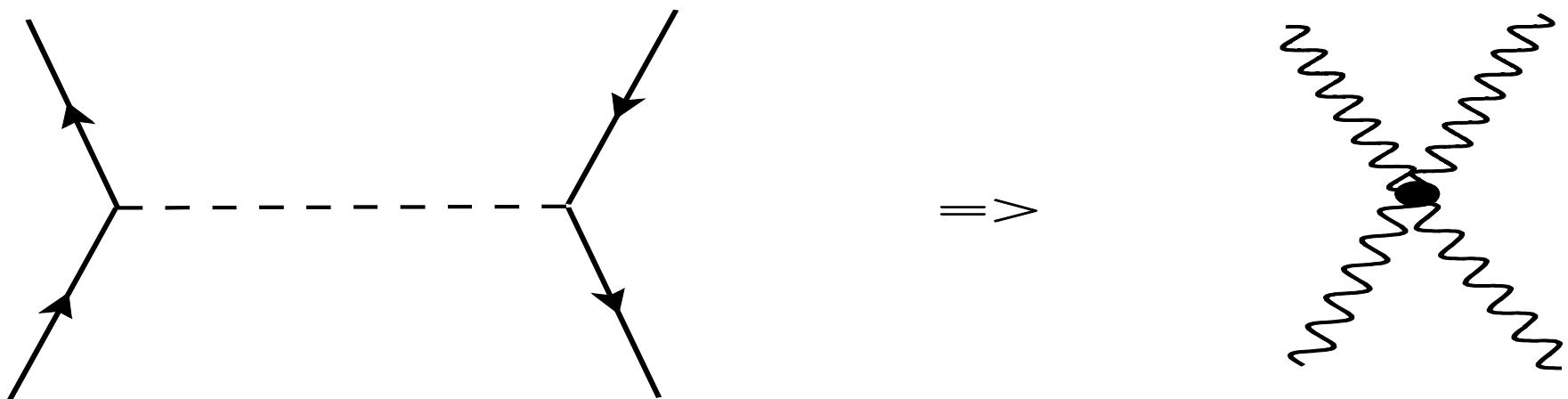
$$\begin{aligned} \mathcal{A}_{(\mu\alpha)(v\gamma)} &= (E_\mu + E_\alpha) \delta_{\alpha\gamma} \delta_{\mu v} \\ &+ \mathcal{V}_{(\mu\alpha)(v\gamma)} \end{aligned}$$

$$|n; \beta\rangle = \sum_{va} \mathbf{C}_{va}^\beta \mathbf{O}_v^\dagger |n-1; \alpha\rangle$$

$$\begin{aligned} \epsilon_p - \epsilon_h &\Rightarrow E_\mu \\ |\text{ph}\rangle &\Rightarrow |\mu\rangle = \sum_{\text{ph}} c^\mu_{\text{ph}} a^\dagger_{\text{p}} a_{\text{h}} | \end{aligned}$$



$$V_{\text{ph}'\text{hp}} \Rightarrow \mathcal{V}_{(\mu v)(\mu' v')}$$



# Generation of n-phonon states

$$\mathbf{A} \mathbf{c} = \hbar \omega \mathbf{c}$$



$$|\lambda\rangle = \sum_{\text{ph}} \mathbf{c}_{\text{ph}}^\lambda \mathbf{a}_p^\dagger \mathbf{a}_h |>$$



$$\mathcal{H}C = E D C$$

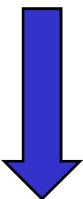


$$|n; \beta\rangle = \sum_{v\alpha} C_{v\alpha}^\beta O_v^\dagger |n-1; \alpha\rangle$$

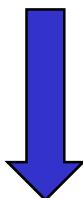
n= 2, 3,.....

# Eigenvalue equations in phonon space

$$\langle \mathbf{n}, \alpha | \mathbf{H} | \Psi_{\nu} \rangle = \mathbf{E}_{\nu} | \Psi_{\nu} \rangle$$



$$\sum_{\mathbf{n}'}^{\beta} [(\mathbf{E}_{\nu} - \mathbf{E}_{\mathbf{n}}^{\alpha}) \delta_{\alpha\beta} \delta_{\mathbf{nn}'} + V^{\alpha\beta}_{(\mathbf{n}, \mathbf{n}')} ] C^{\nu}_{\mathbf{n}'\beta} = 0$$

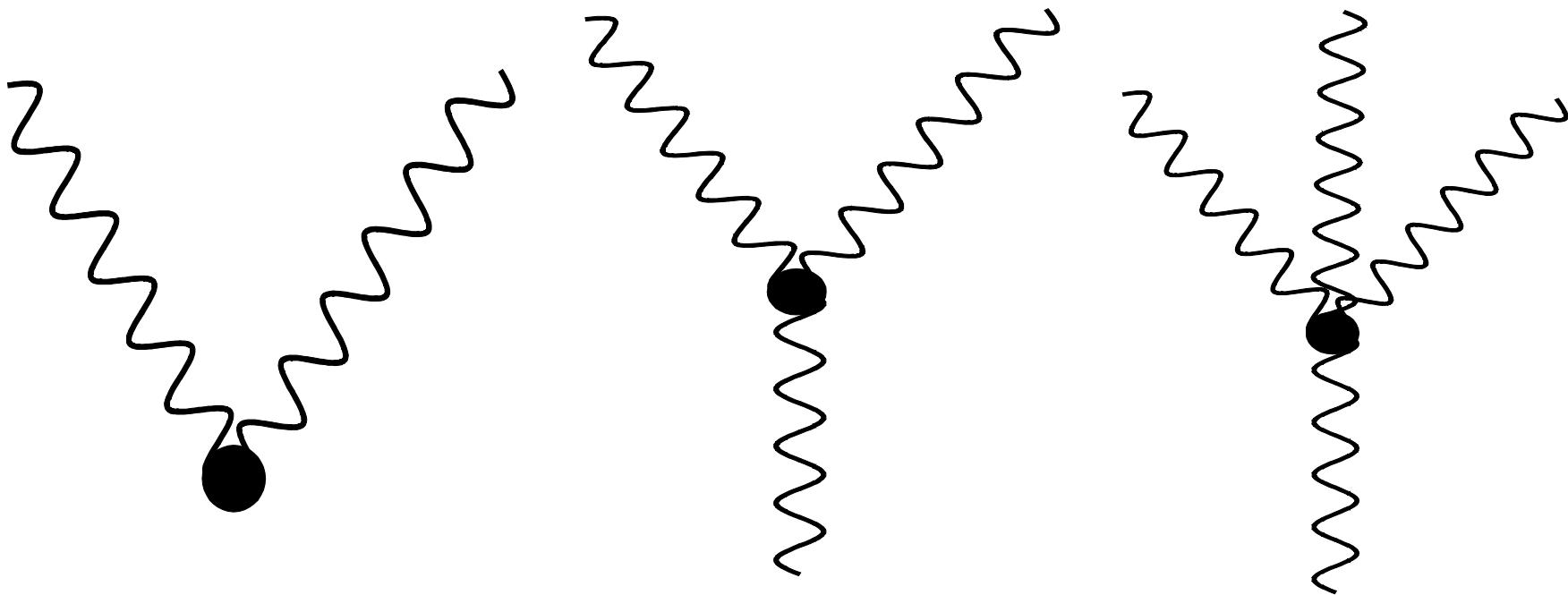


$$| \Psi_{\nu} \rangle = \sum_{\mathbf{n}\alpha} C^{\nu}_{\mathbf{n}\alpha} | \mathbf{n}, \alpha \rangle$$

# Hamiltonian in the $|n-\alpha\rangle$ basis ( $n=0,1,2,\dots$ )

$E^{(0)}$	$\langle 0\text{ph} H 1\text{ph}\rangle$	$\langle 0\text{ph} H 2\text{ph}\rangle$	0
$\langle 1\text{ph} H 0\text{ph}\rangle$	$E^{(1)}_1$ $E^{(1)}_2$ $O$ $O$ $E^{(1)}_k$	$\langle 1\text{ph} H 2\text{ph}\rangle$	$\langle 1\text{ph} H 3\text{ph}\rangle$
$\langle 2\text{ph} H 0\text{ph}\rangle$	$\langle 2\text{ph} H 1\text{ph}\rangle$	$E^{(2)}_1$ $E^{(2)}_2$ $O$ $O$ $E^{(2)}_k$	$\langle 2\text{ph} H 3\text{ph}\rangle$
0	$\langle 3\text{ph} H 1\text{ph}\rangle$	$\langle 3\text{ph} H 2\text{ph}\rangle$	$E^{(3)}_1$ $E^{(3)}_2$ $O$ $O$ $E^{(3)}_k$

Diagrammatic:  $V_{nn'}$ ,

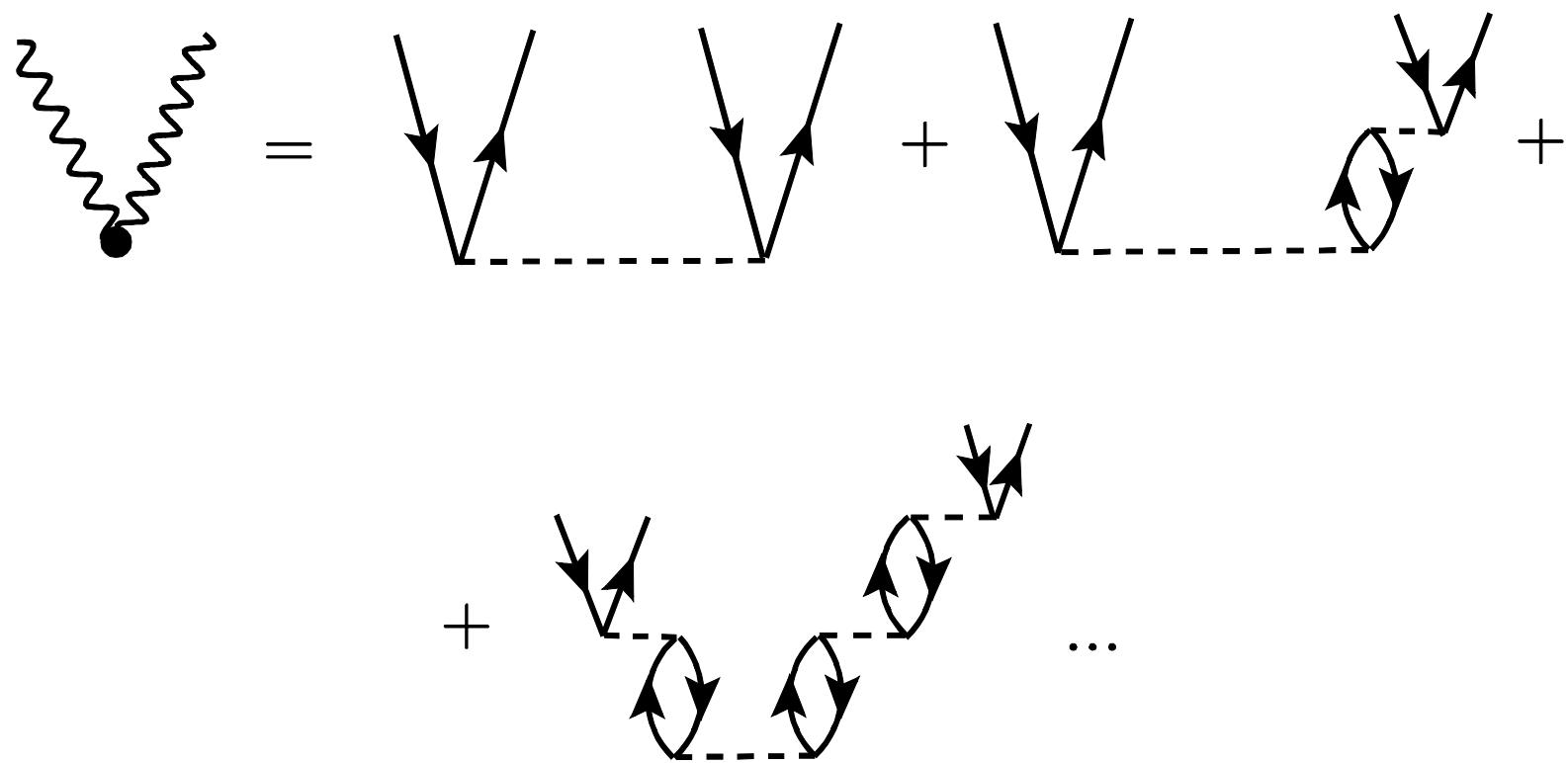


$$\boxed{<2|V|0>}$$

$$\boxed{<2|V|1>}$$

$$\boxed{<3|V|1>}$$

# Ground State Correlations (**real!**)



# EMPM : Numerical implementation

Hamiltonian

$$H = H_0 + V = \sum_i h_i + G_{\text{bare}} \quad (V_{\text{BonnA}} \Rightarrow G_{\text{bare}})$$

$$h = t + h_{\text{Nils}}$$

$$h = t + h_{\text{HF}}$$

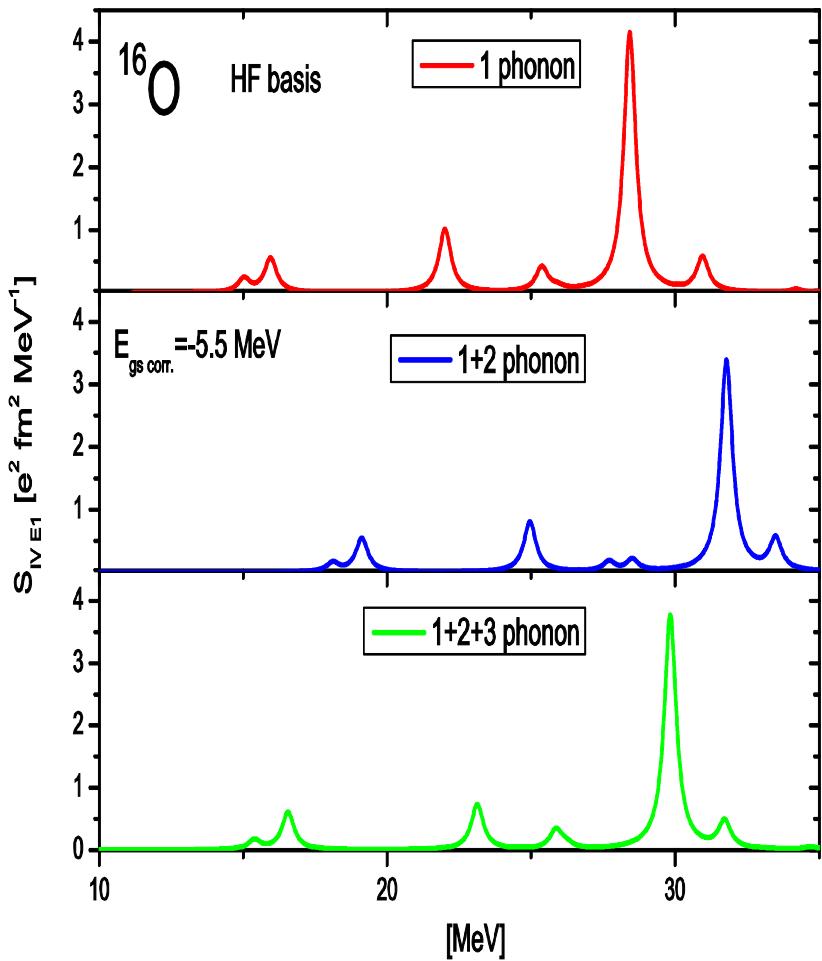
Application to O isotopes

Phonon space: up to n=3

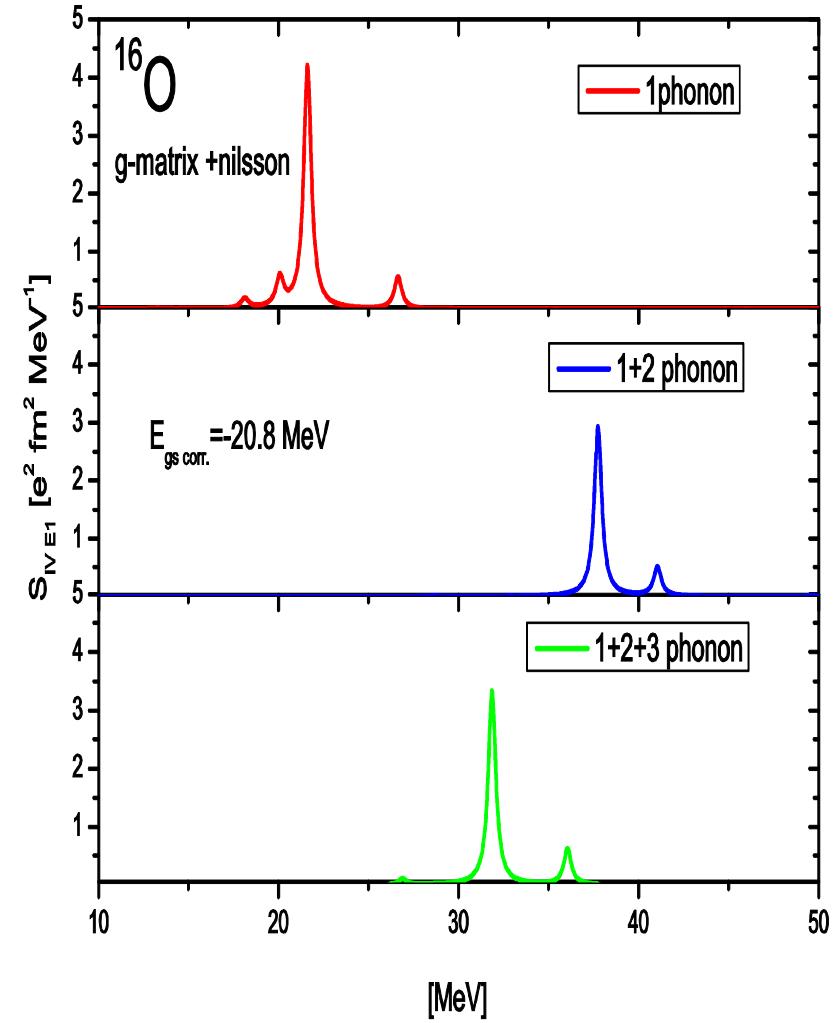
SM space

1. p-h of  $1\hbar\omega$
2. p-h of  $3\hbar\omega$  (with truncation of the n=3 phonon space)

# $^{16}\text{O}$ : E1 Strength Function



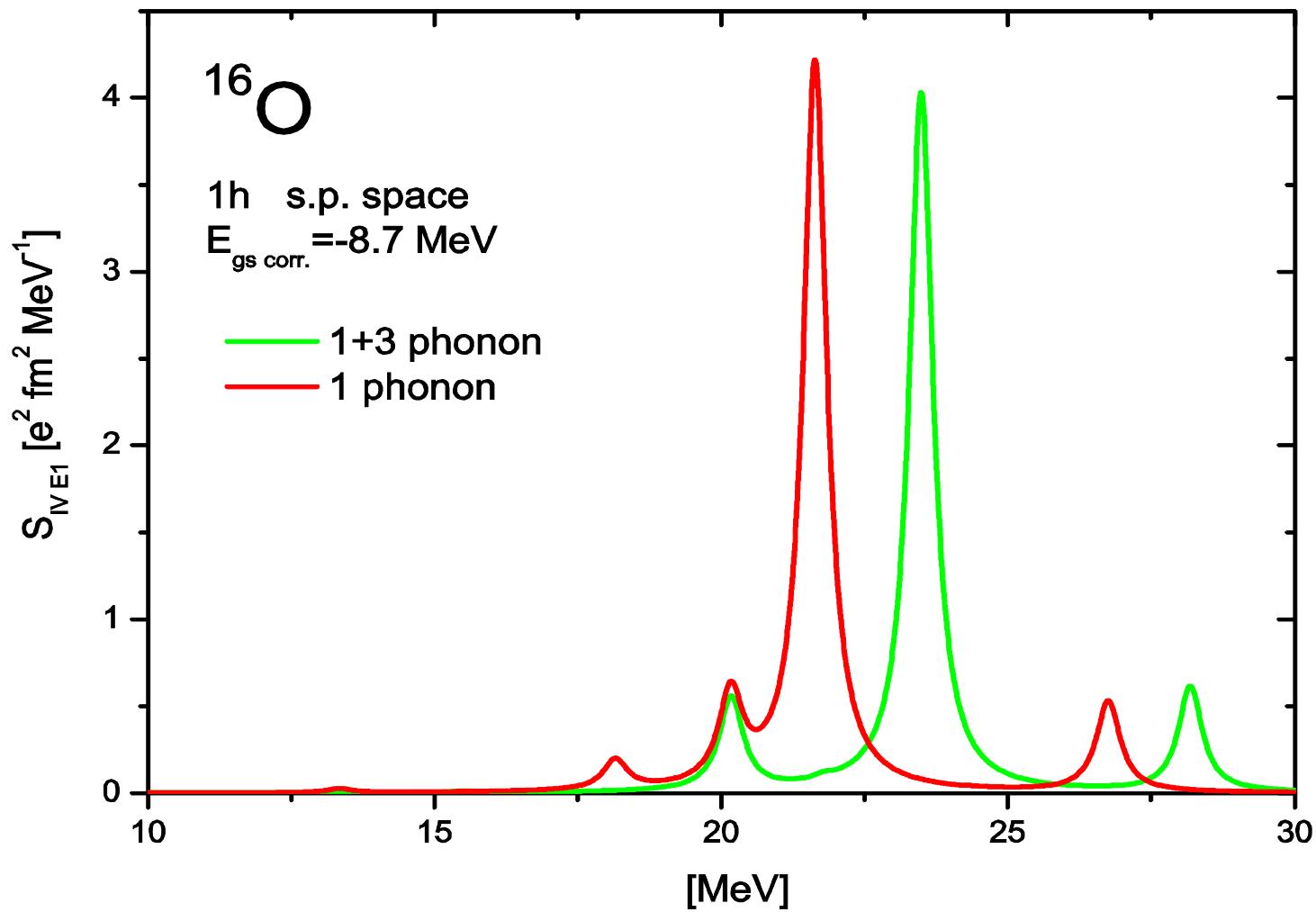
HF



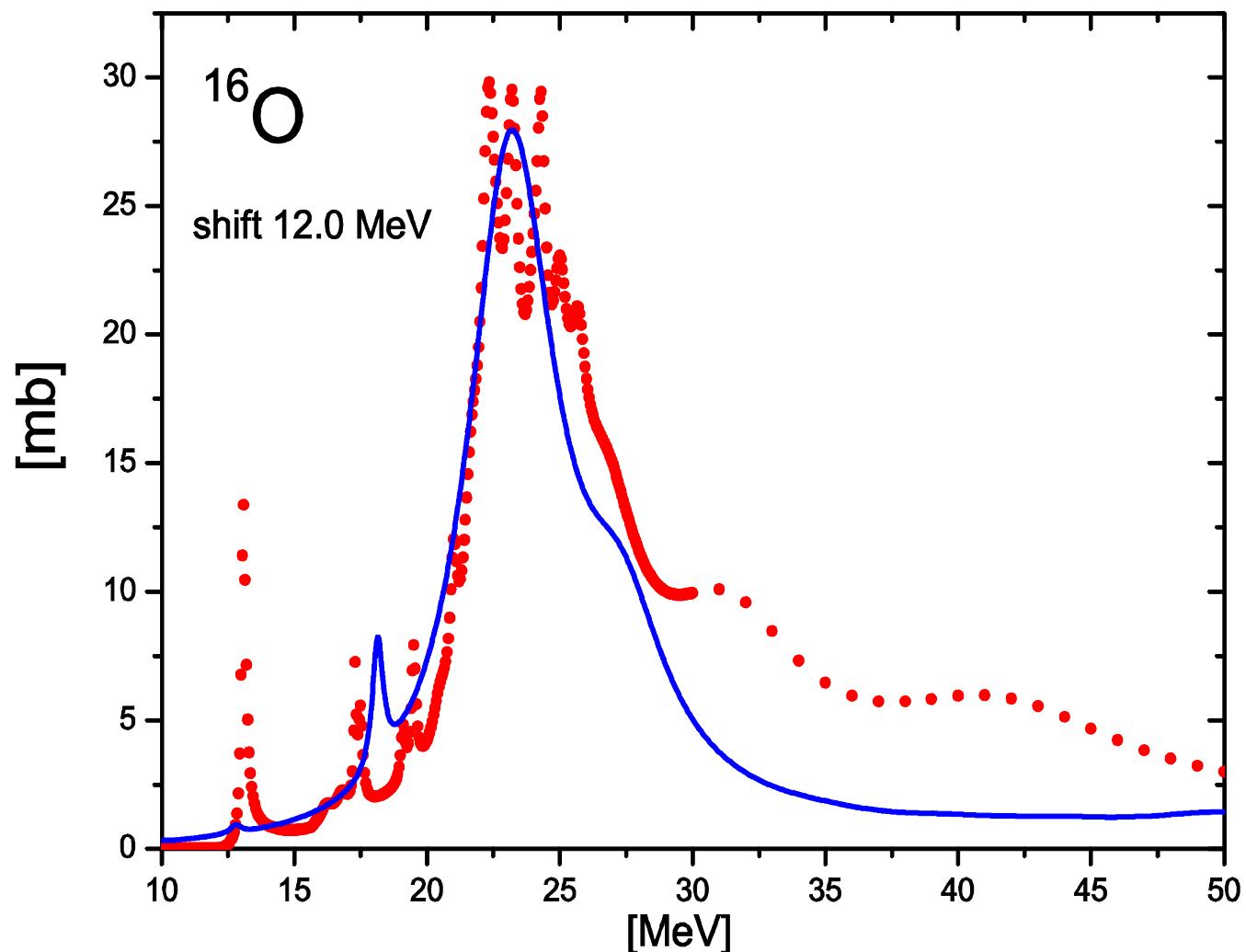
Nilsson

# GDR

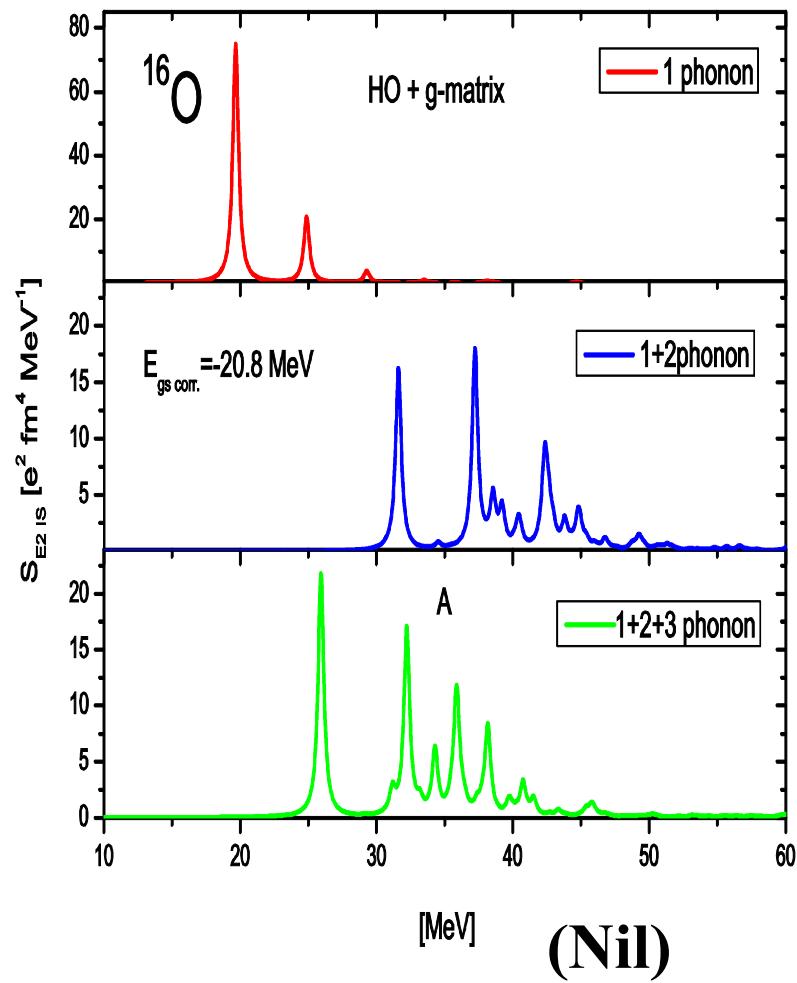
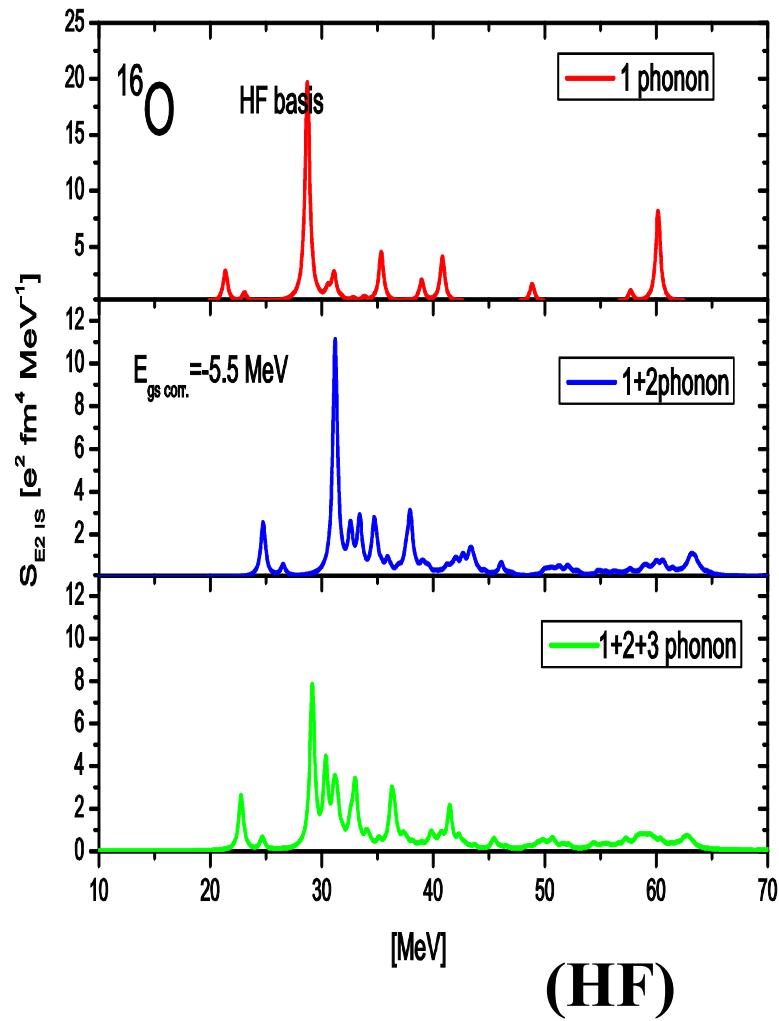
# (1 $\hbar\omega$ )



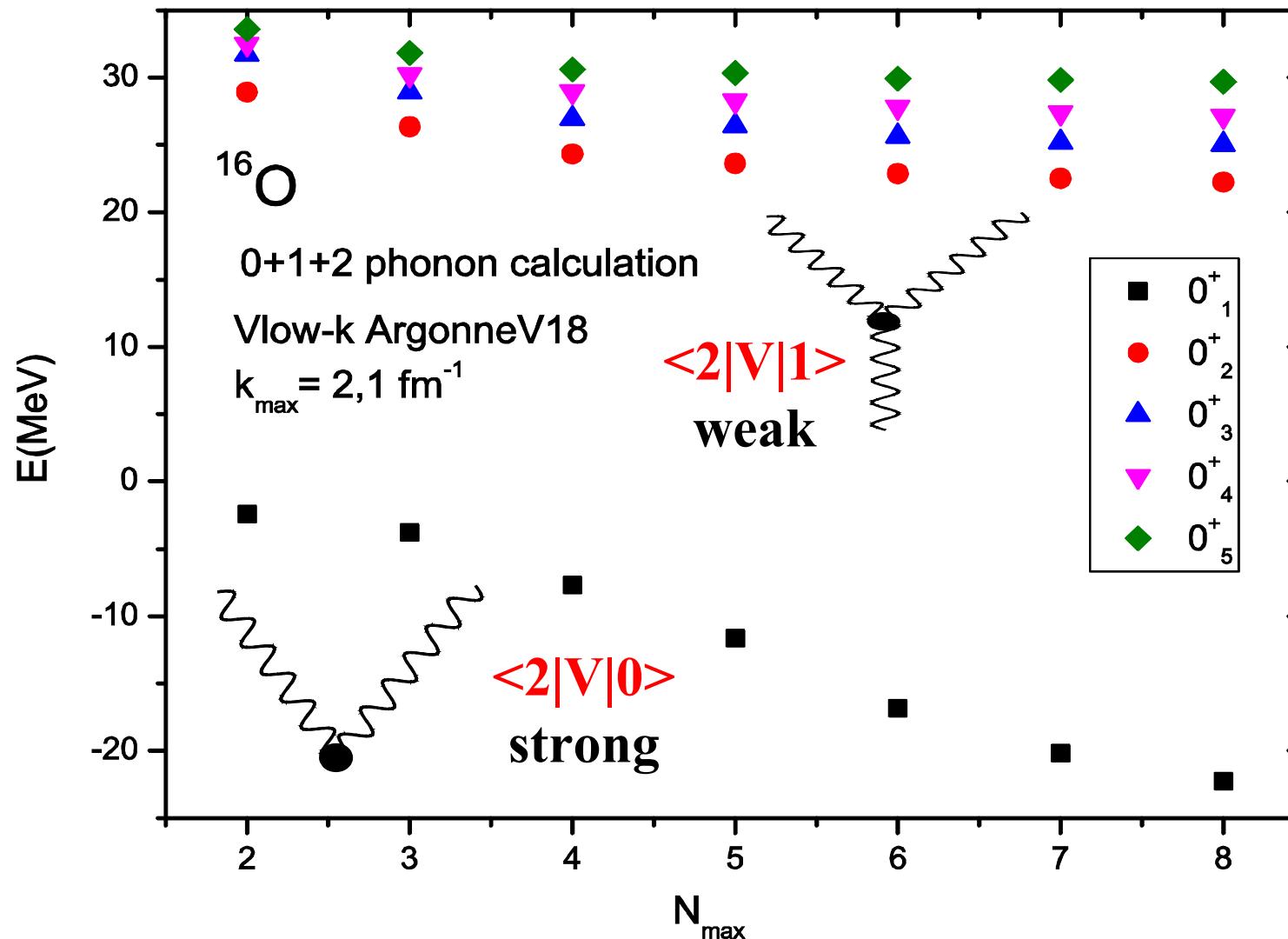
# GDR: Cross Section



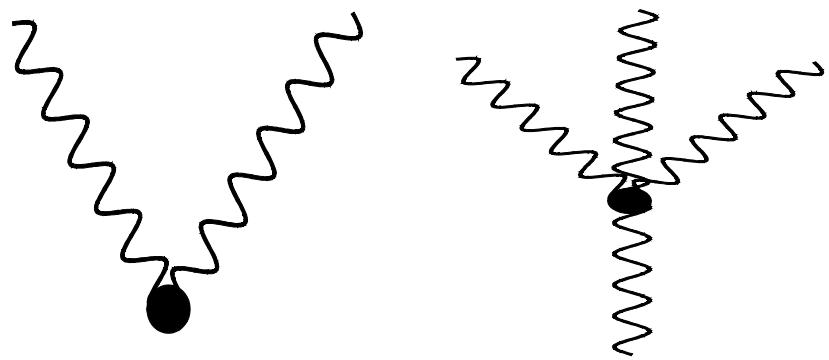
# E2 Strength Function



# Ground state

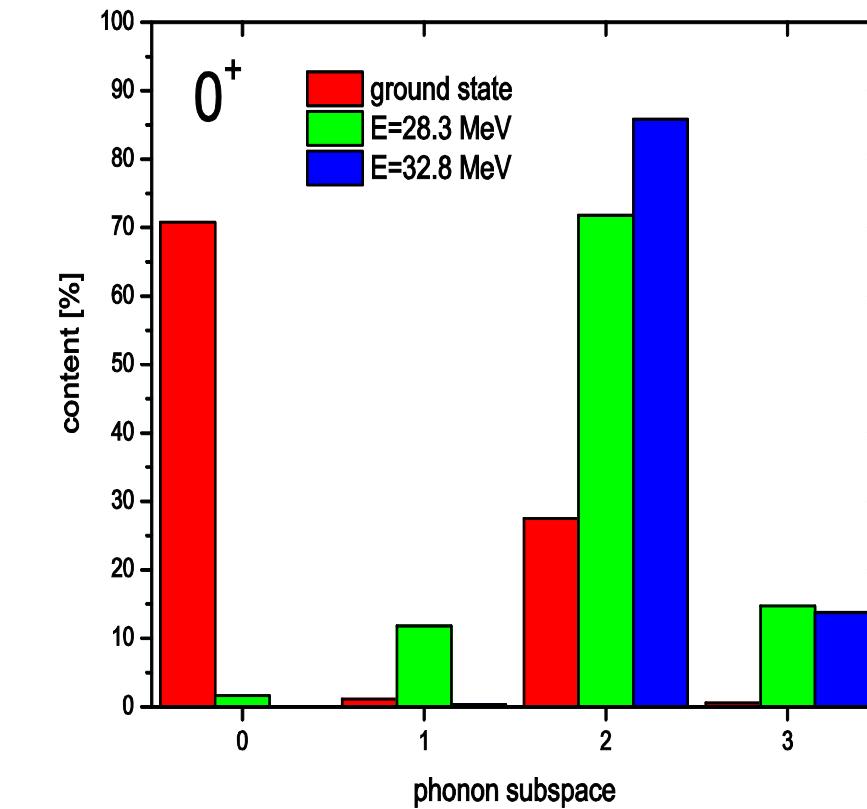


# 3-phonons needed but not enough



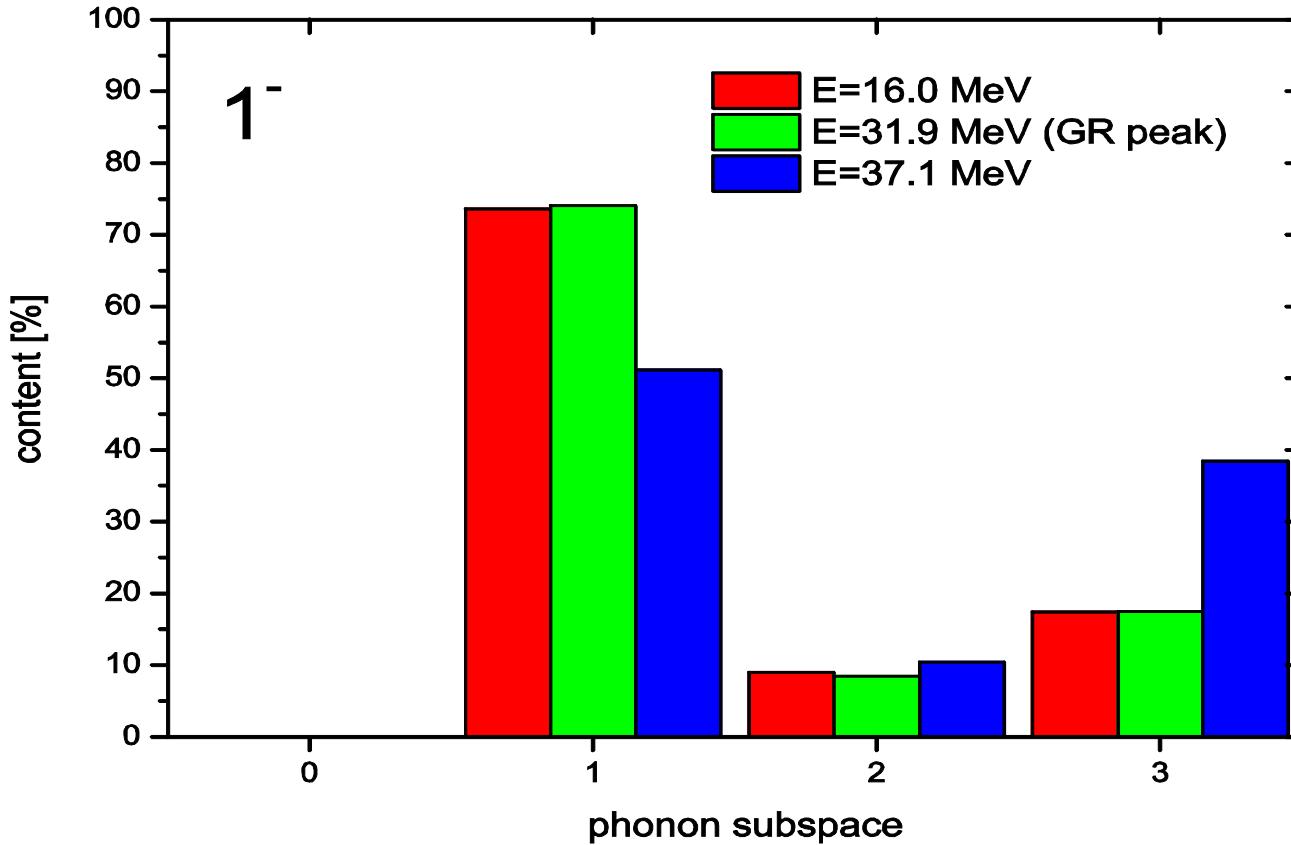
$$\langle n | V | n-2 \rangle$$

strong



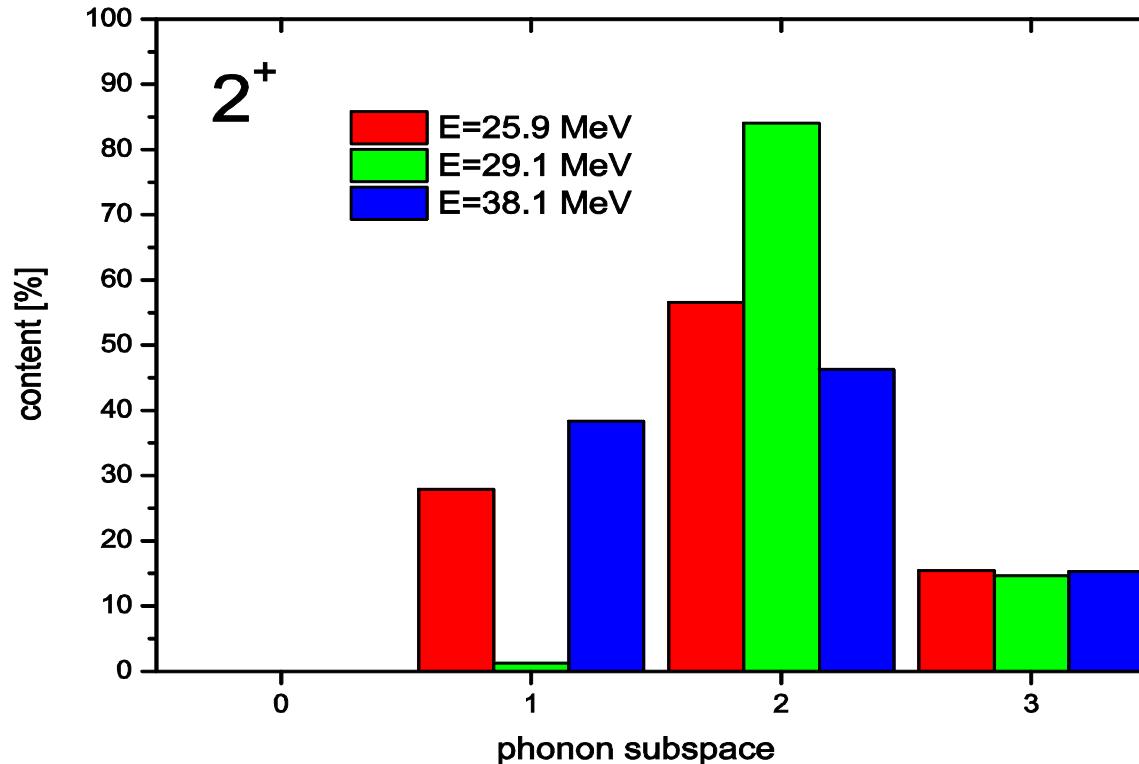
# 4-phonons needed

# Phonon content: $1^-$



3-phonon maybe enough but large p-h space

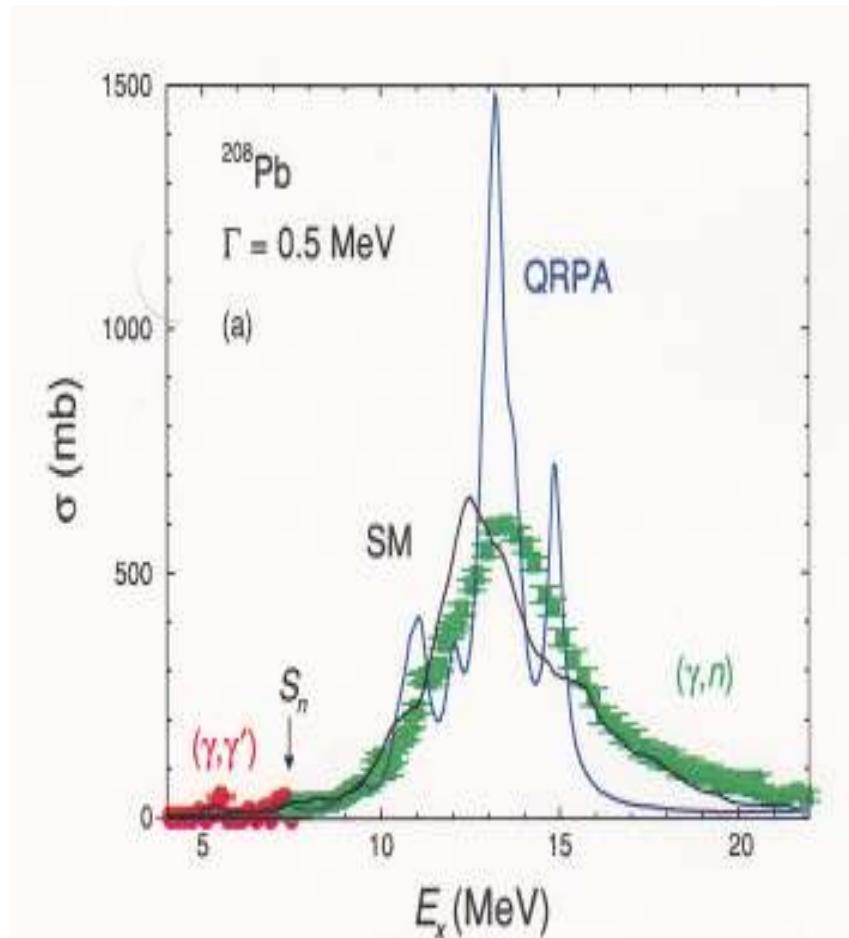
# Phonon content: $2^+$



**4-phonons** needed also for  $2^+$

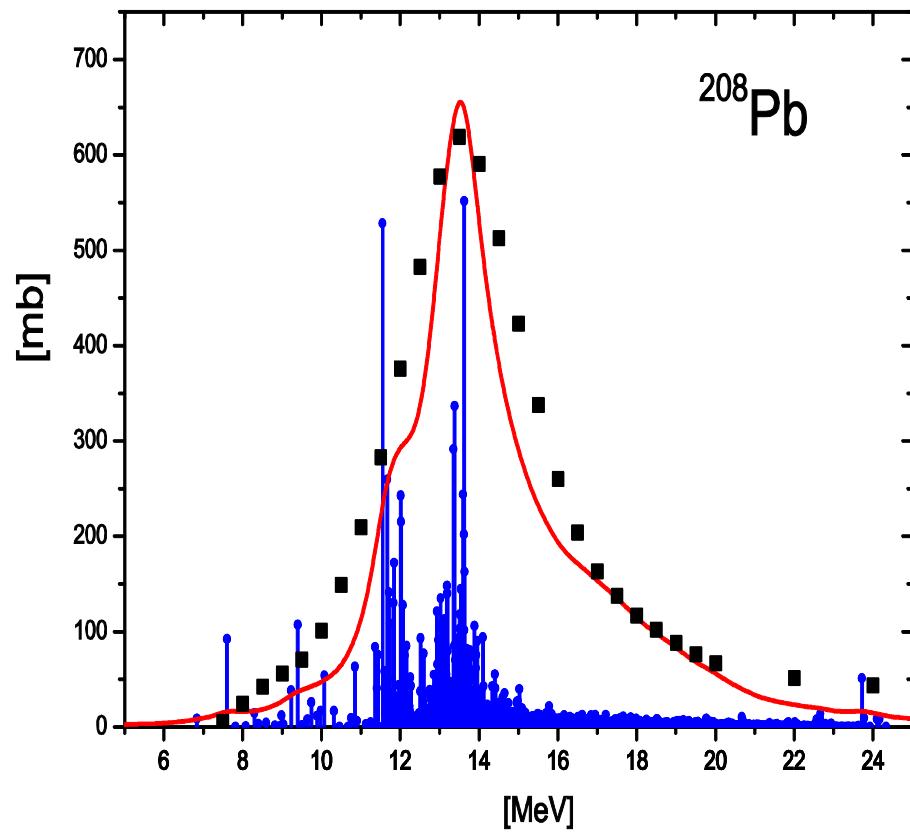
# Application to heavy nuclei (easy up to **two-phonons**)

## GDR in $^{208}\text{Pb}$



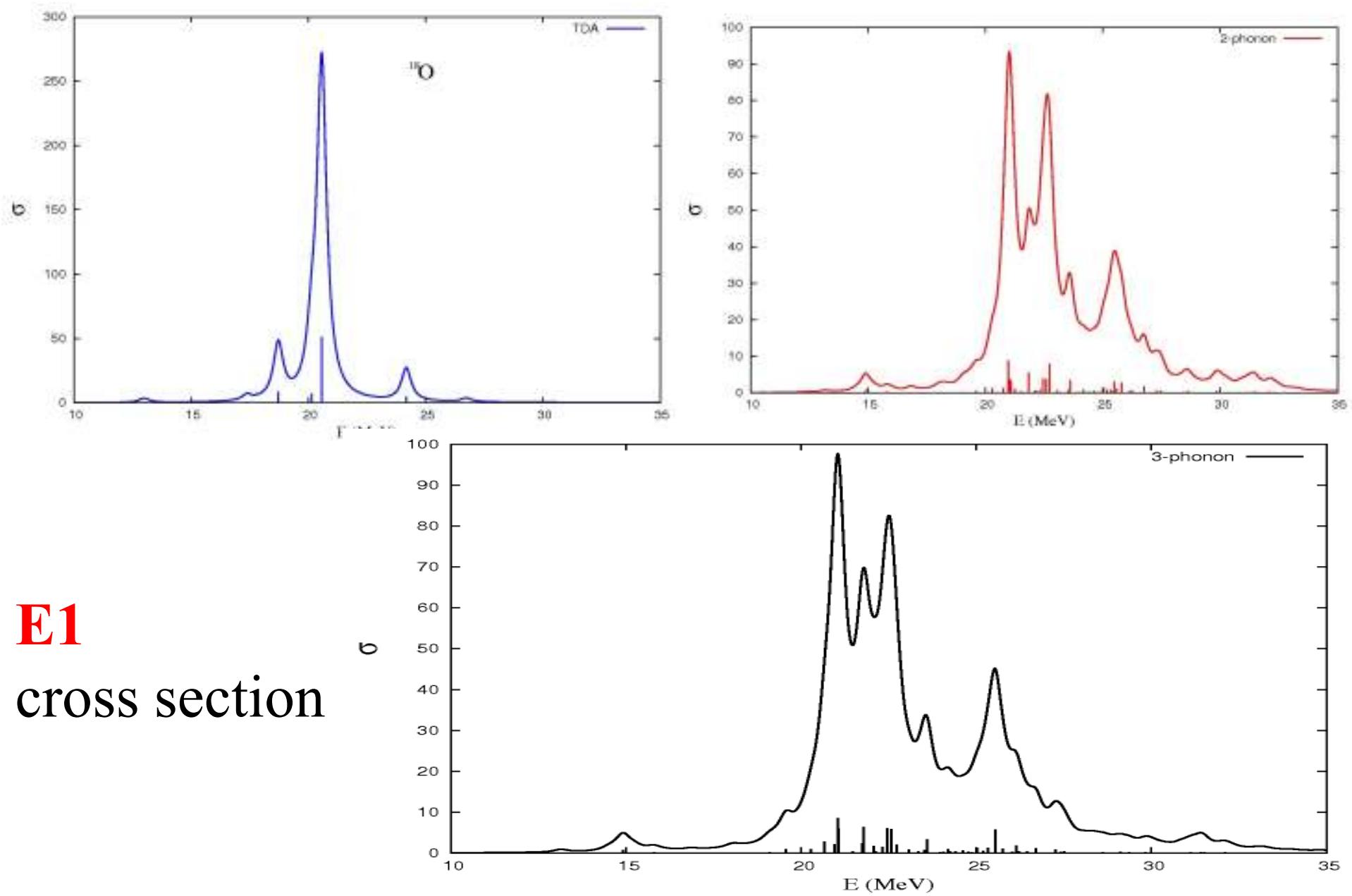
SM

(R. Schwengner..A. Brown PRC81(2010))

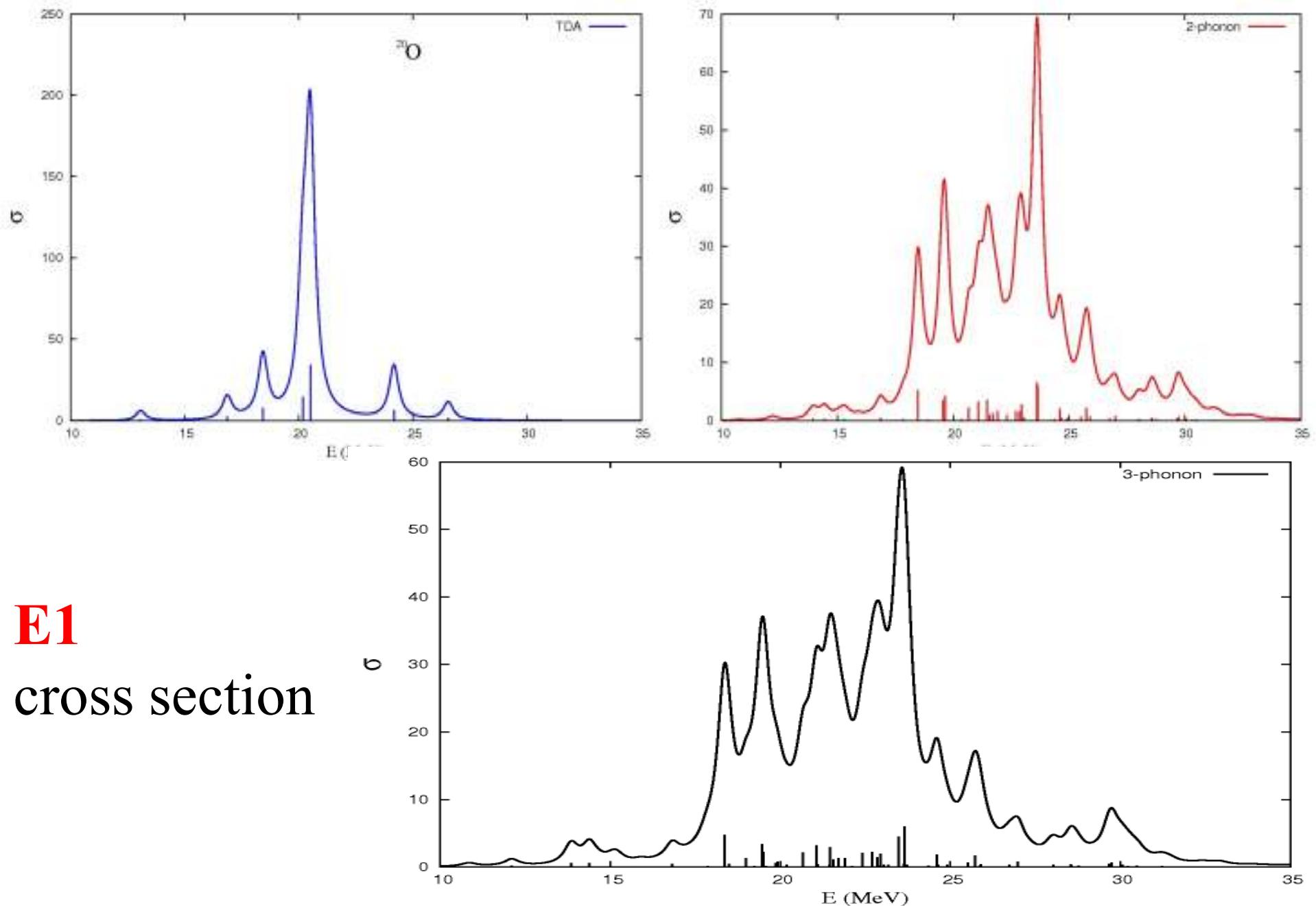


EMPM

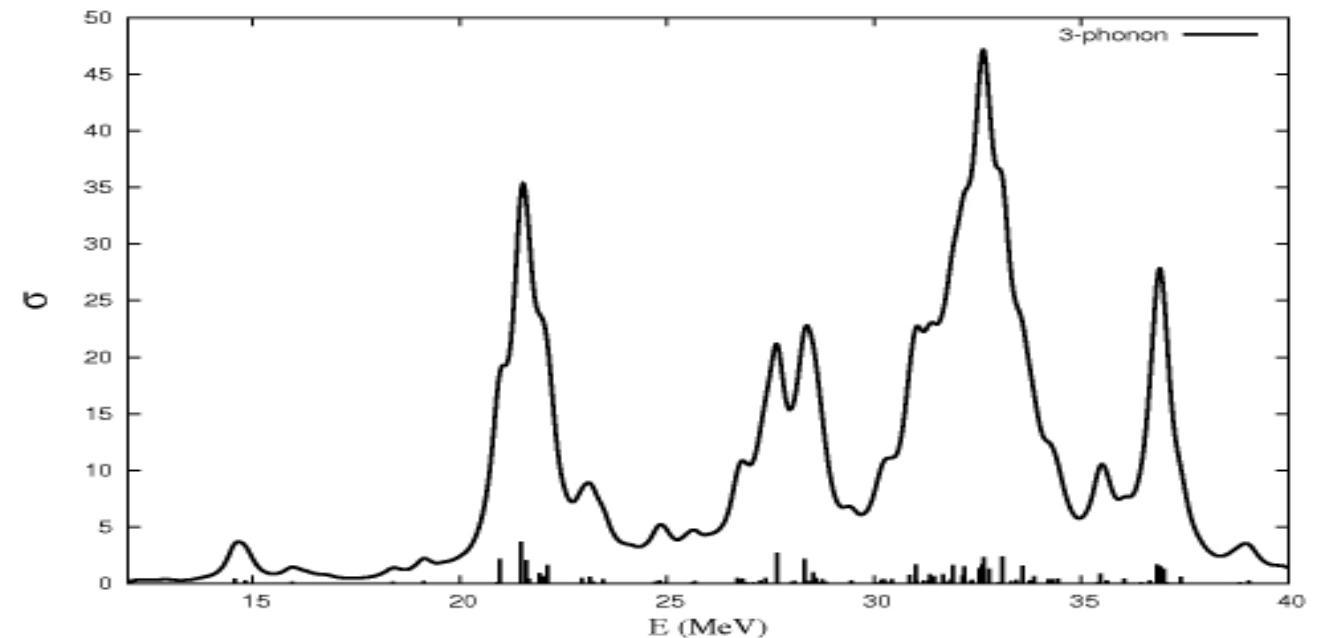
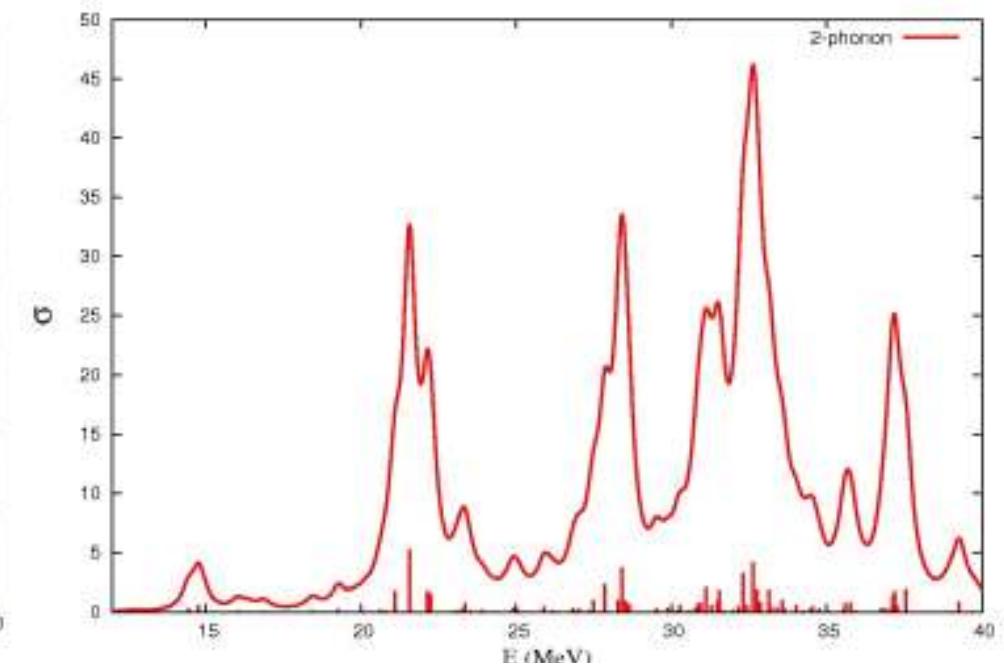
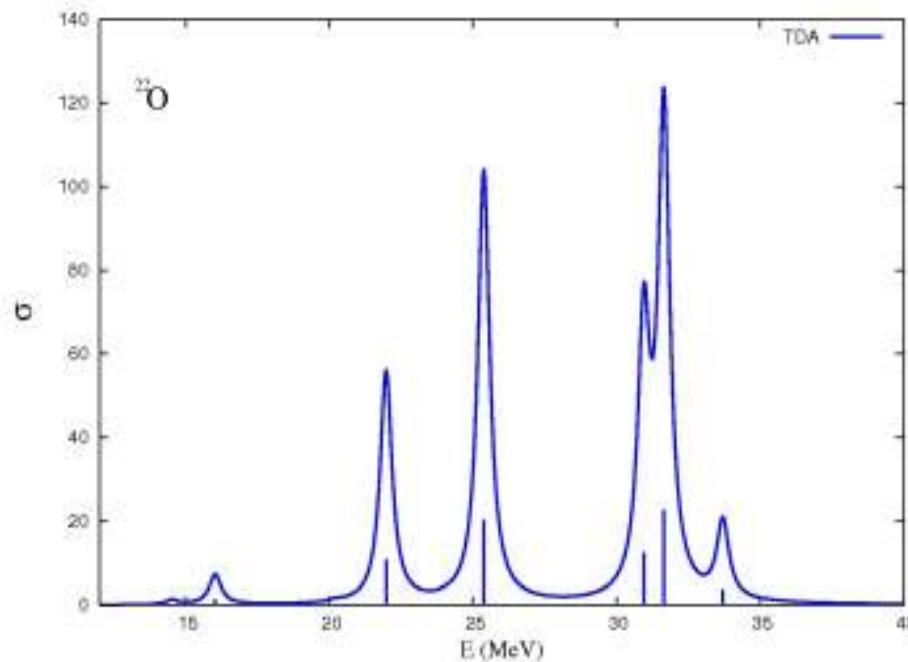
# Open shell nuclei: $^{18}\text{O}$



# Open shell nuclei: $^{20}\text{O}$

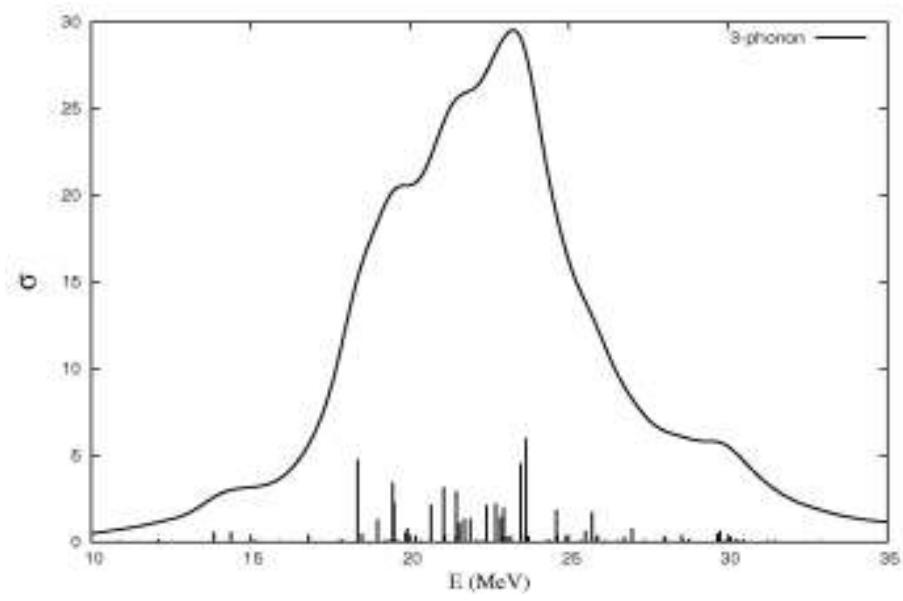
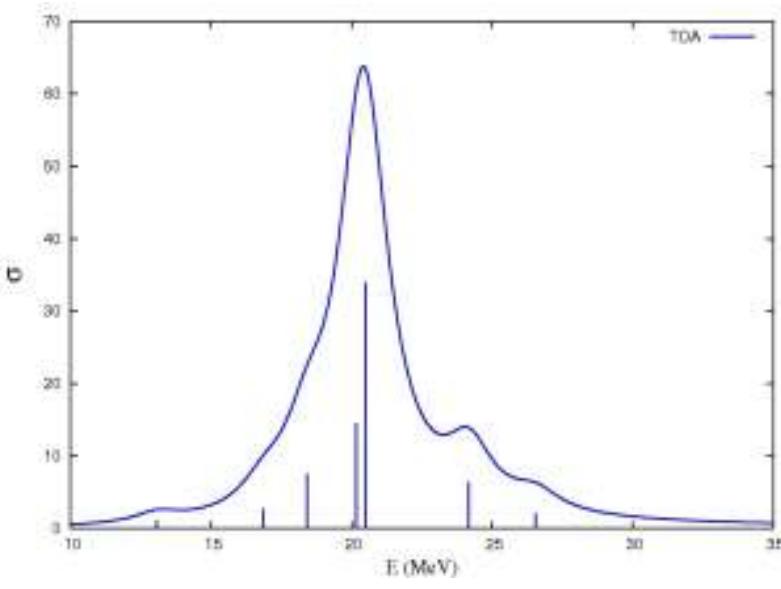
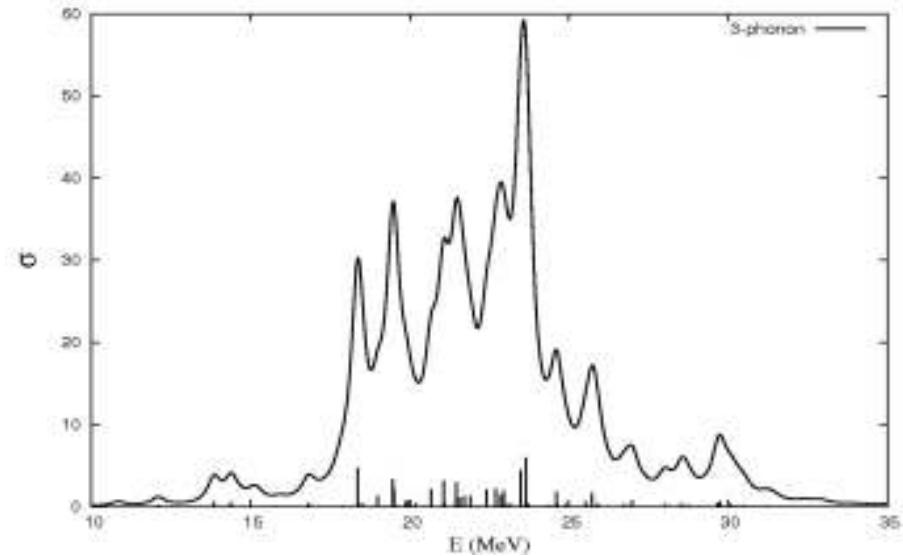
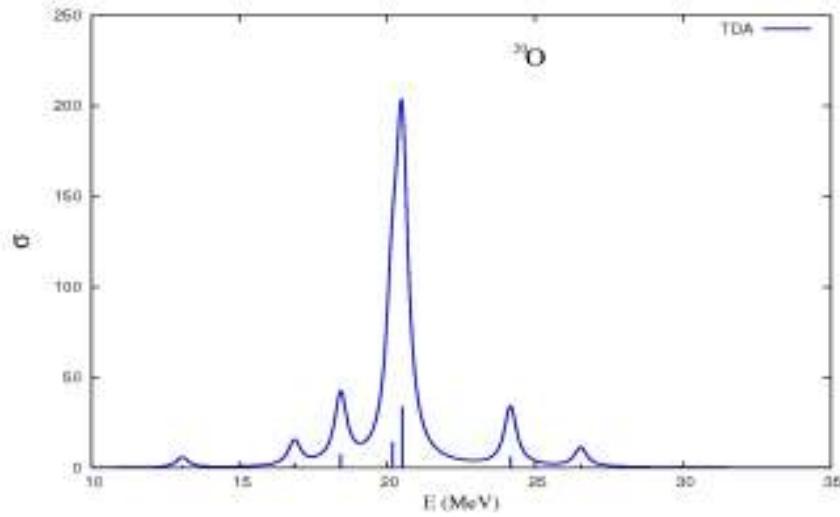


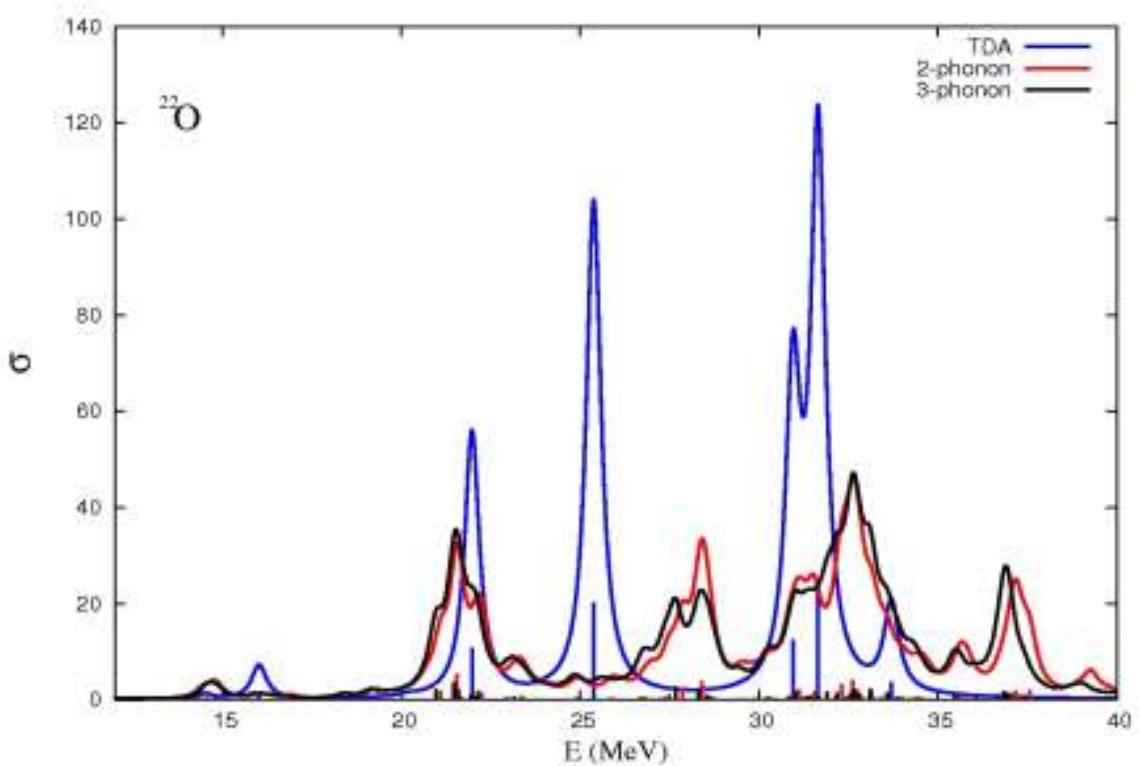
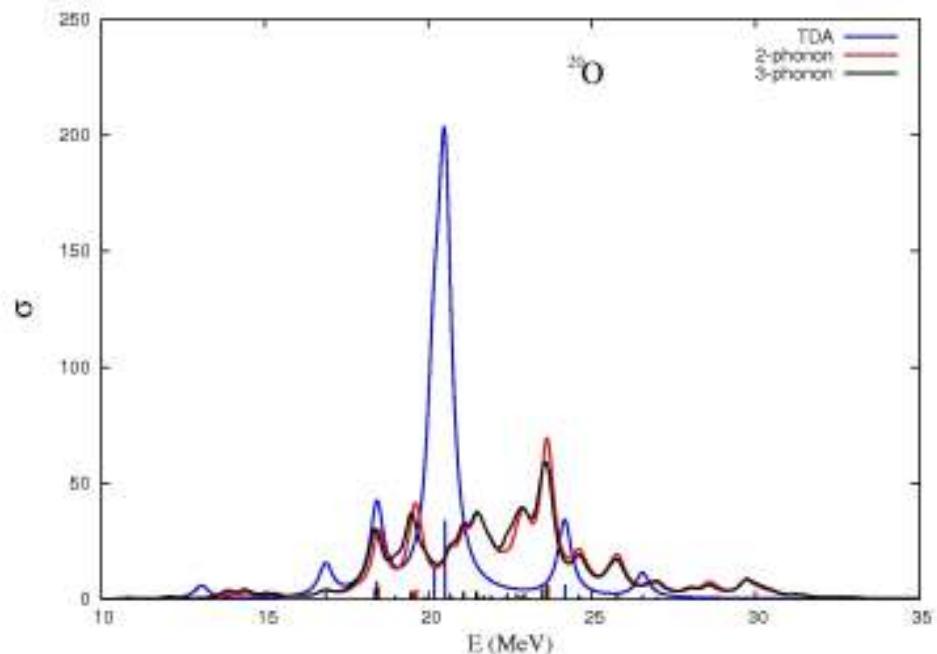
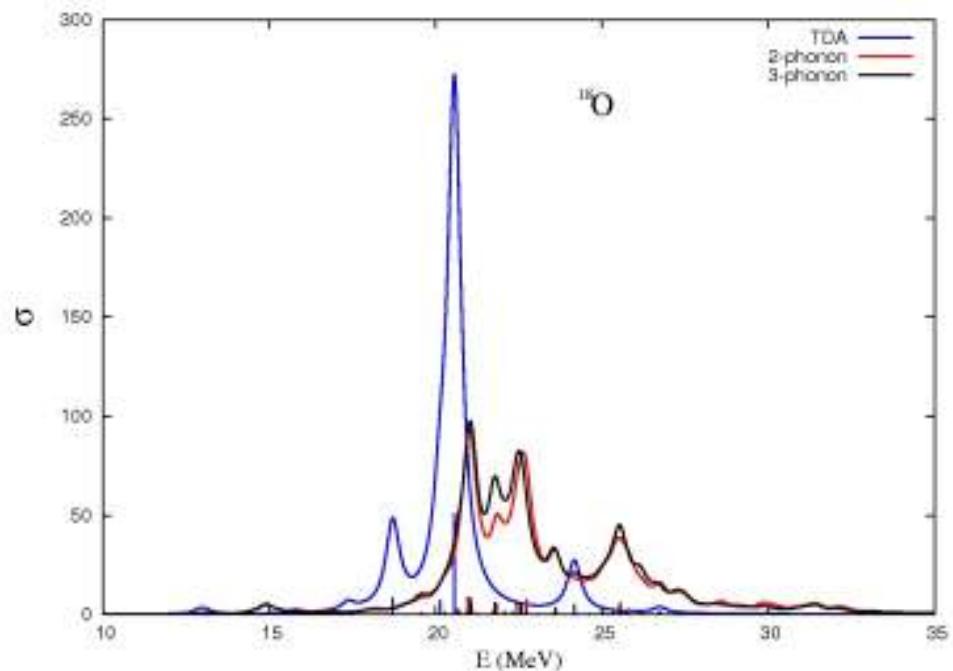
# Open shell nuclei: $^{22}\text{O}$



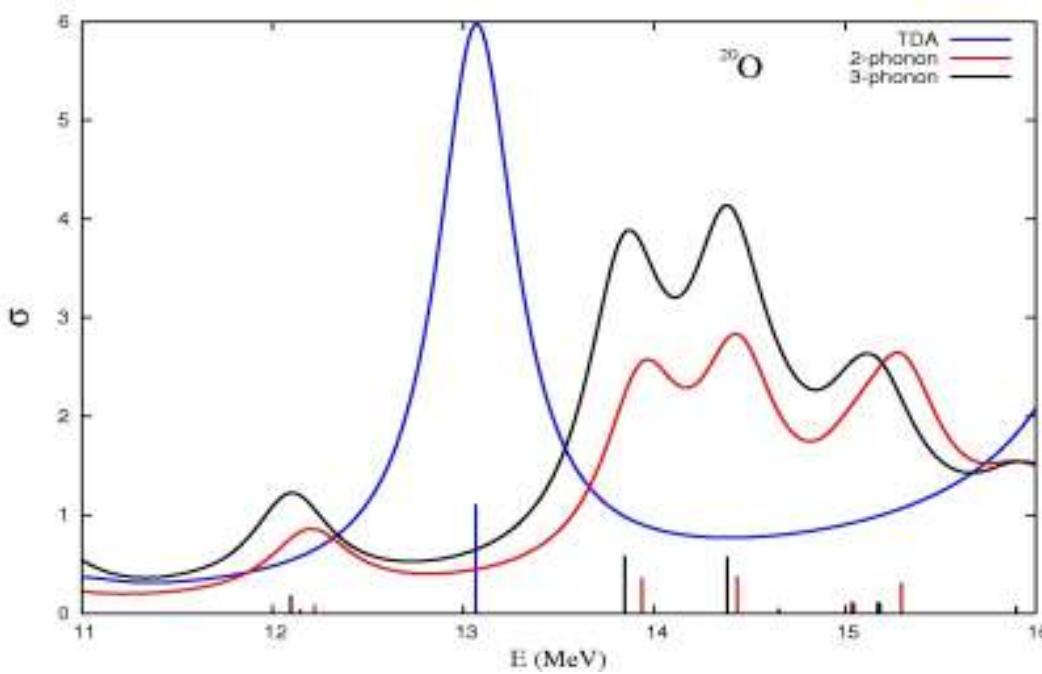
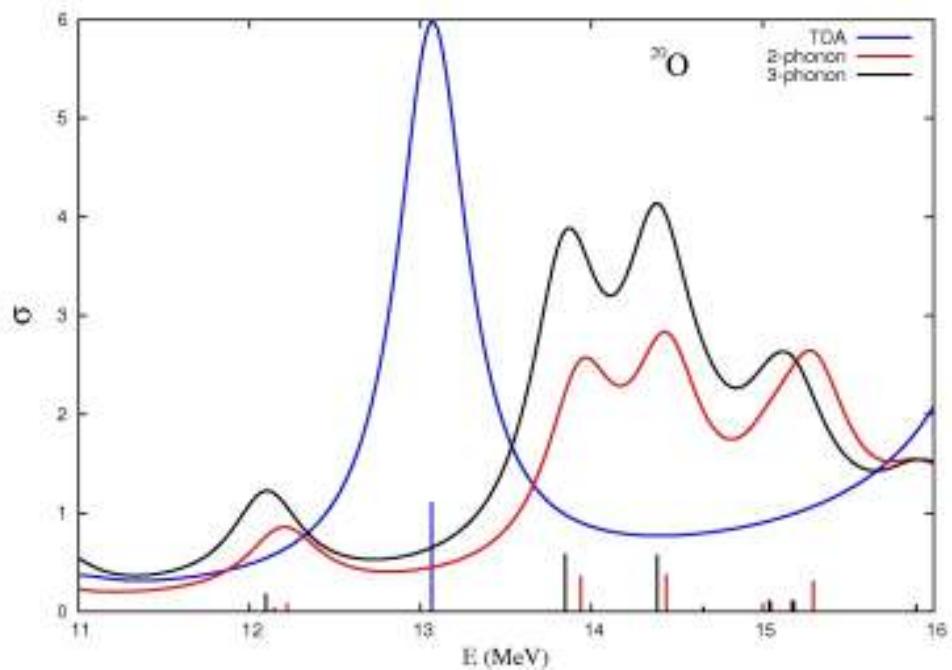
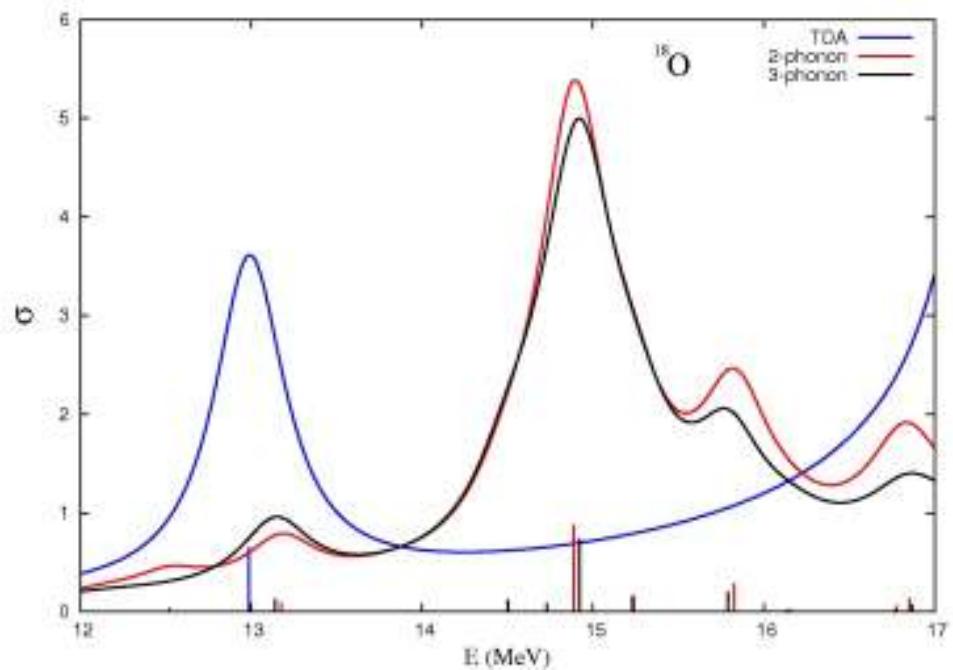
E1  
cross section

# Open shell nuclei: $^{20}\text{O}$





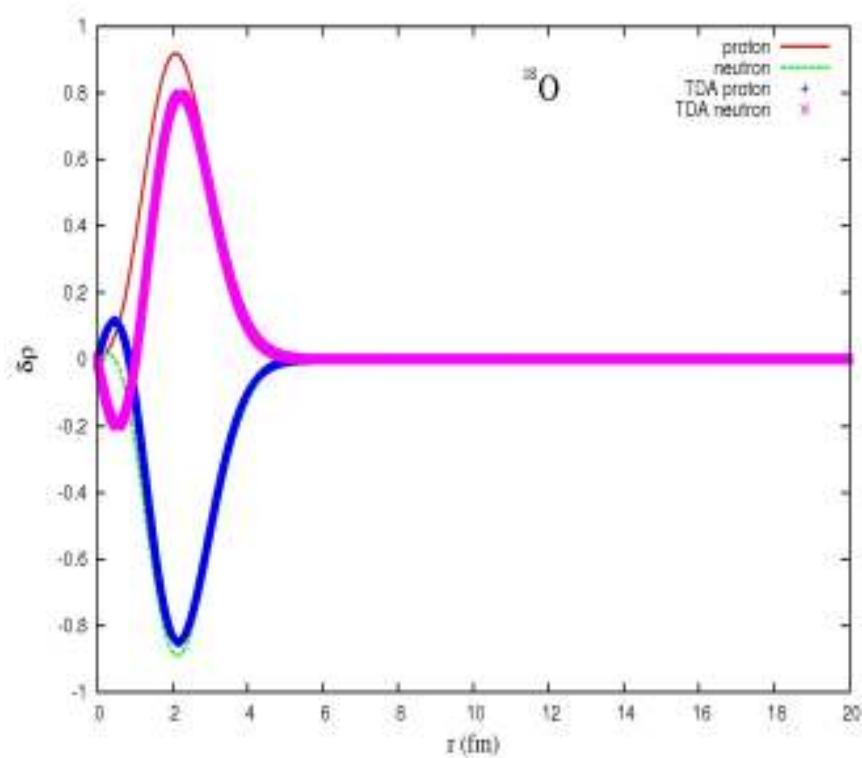
Overview  
E1 cross section  
in  $^{18-22}\text{O}$



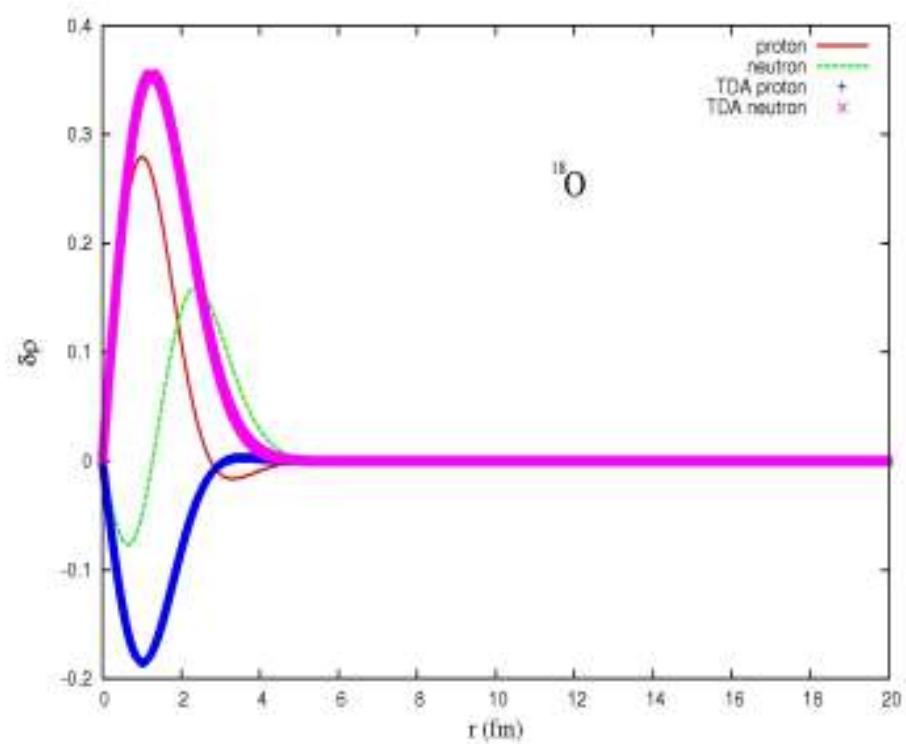
Pygmy

# E1 transition density: $^{18}\text{O}$

GDR

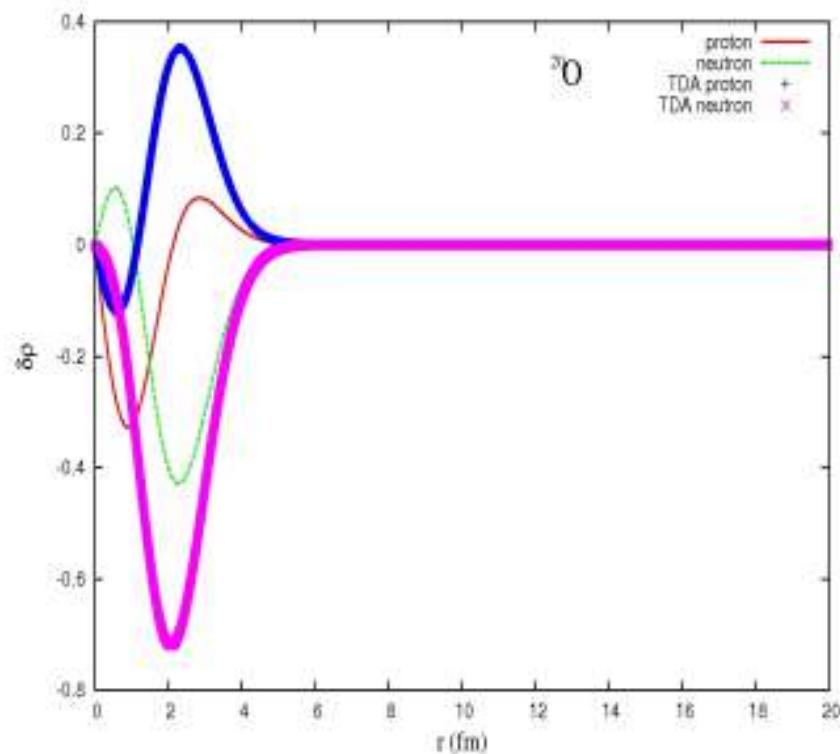


Pygmy

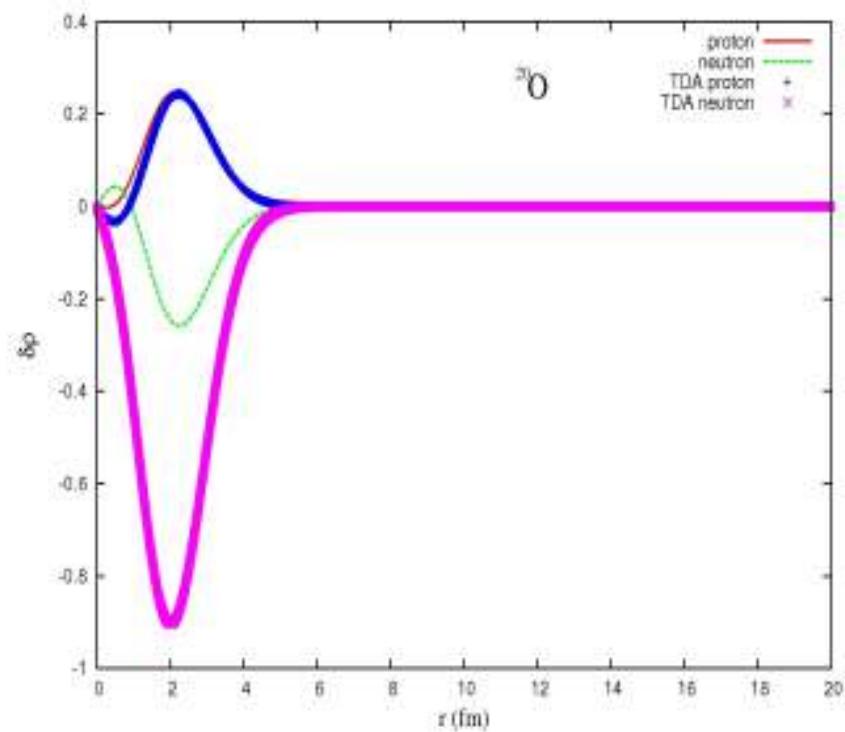


# E1 transition density: $^{20}\text{O}$

GDR

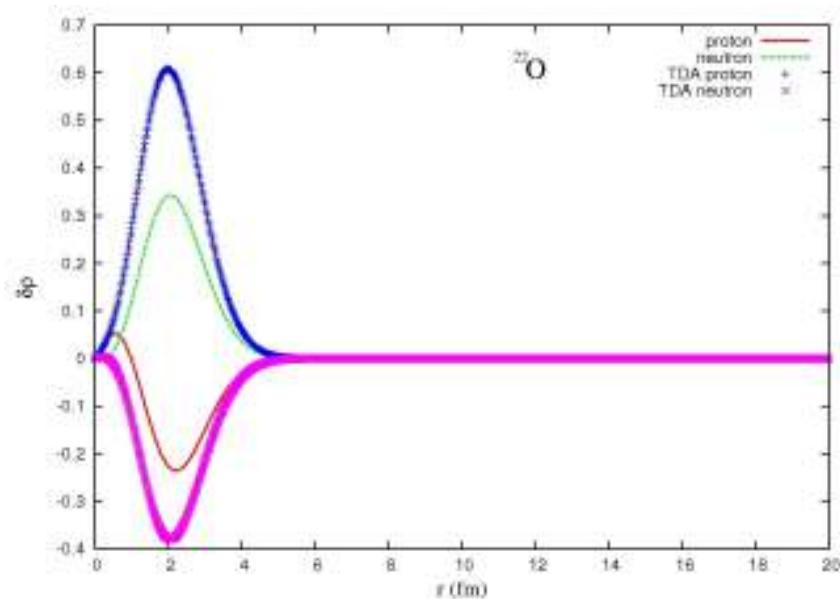


Pygmy

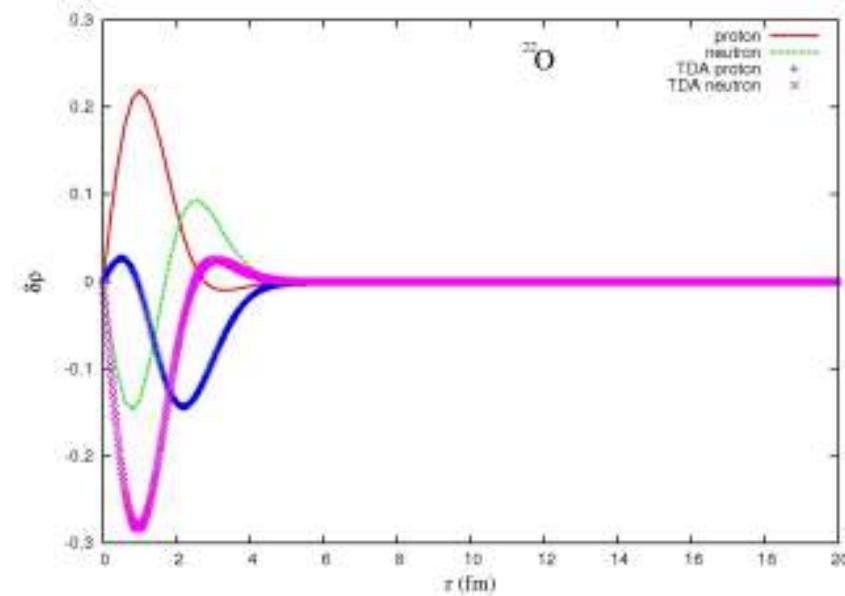


# E1 transition density: $^{22}\text{O}$

GDR



Pygmy



# Concluding remarks

- The **EMPM** gives an **exact formulation** of the **Nuclear Eigenvalue Problem** in a **microscopic phonon** basis for an **arbitrary** number **n** of phonons

- Actual implementation :
  - **n** arbitrarily large in **small** configuration space
  - **n=3** in fairly large **configuration** space (but **n=4** is in progress)
  - **n=2** in very large configuration space

Final goal:

generate **TDA phonons** in **very large** configuration **space** and **select** only the **relevant** ones so as to allow a

**Truncation** of the **phonon space**

**THANK YOU**

# $^{16}\text{O}$ as theoretical lab

## Structure of $^{16}\text{O}$ : A theoretical challenge

Pioneering work: First excited  $0^+$  as deformed 4p-4h excitations

G. E. Brown, A. M. Green, Nucl. Phys. 75, 401 (1966)

(TDA) IBM (includes up to 4 TDA Bosons)

H. Feshbach and F. Iachello, Phys. Lett. B 45, 7 (1973); Ann. Phys. 84, 211 (1974)

SM up to 4p-4h and 4  $\hbar\omega$

W.C. Haxton and C. J. Johnson, PRL 65, 1325 (1990)

E.K. Warbutton, B.A. Brown, D.J. Millener, Phys. Lett. B293, 7(1992)

No-core SM (NCSM) Huge space!!!

Symplectic No-core SM (SpNCSM) a promising tool for cutting the SM space

T. Dytrych, K.D. Sviratcheva, C. Bahri, J. P. Draayer, and J.P. Vary, PRL 98, 162503 (2007)

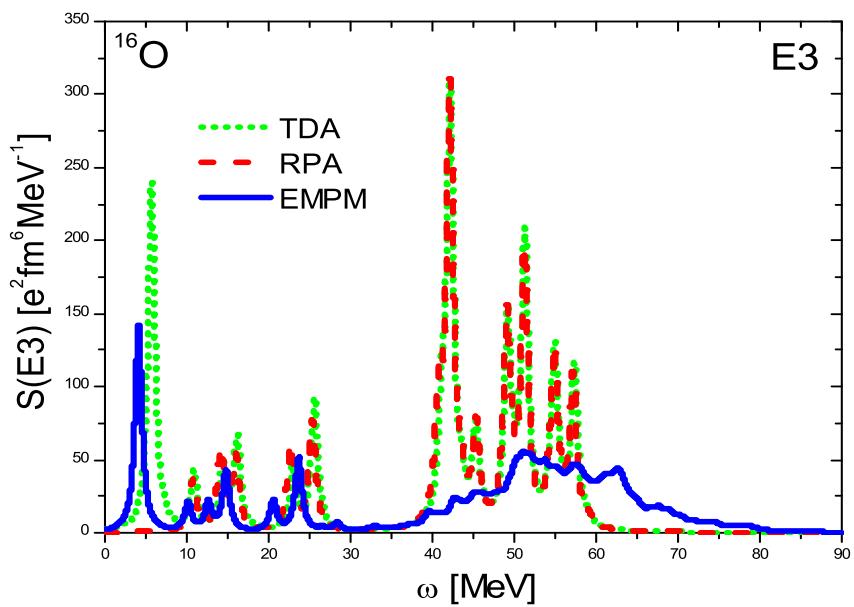
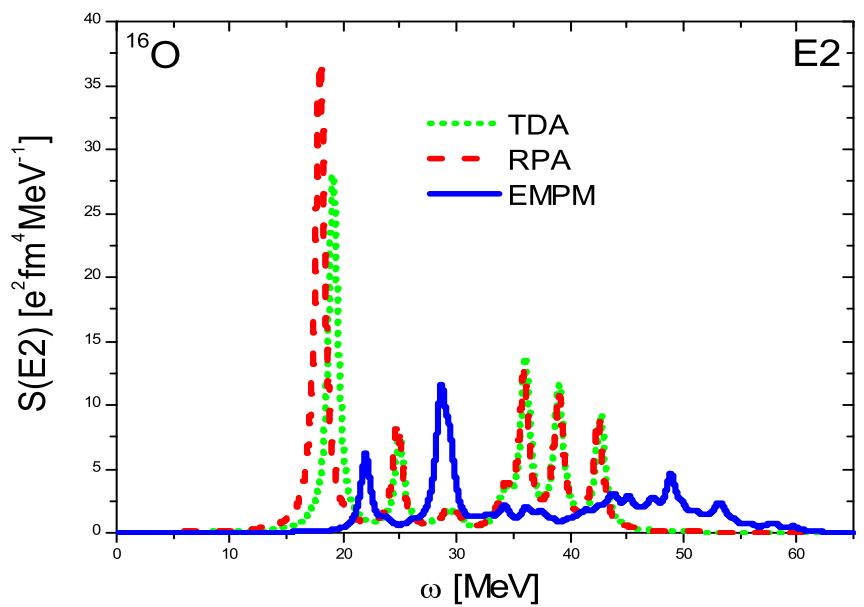
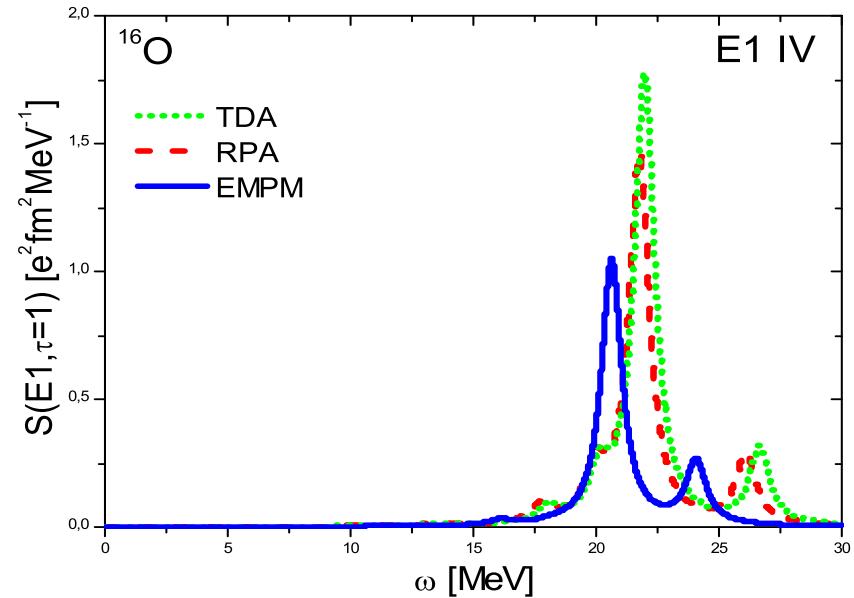
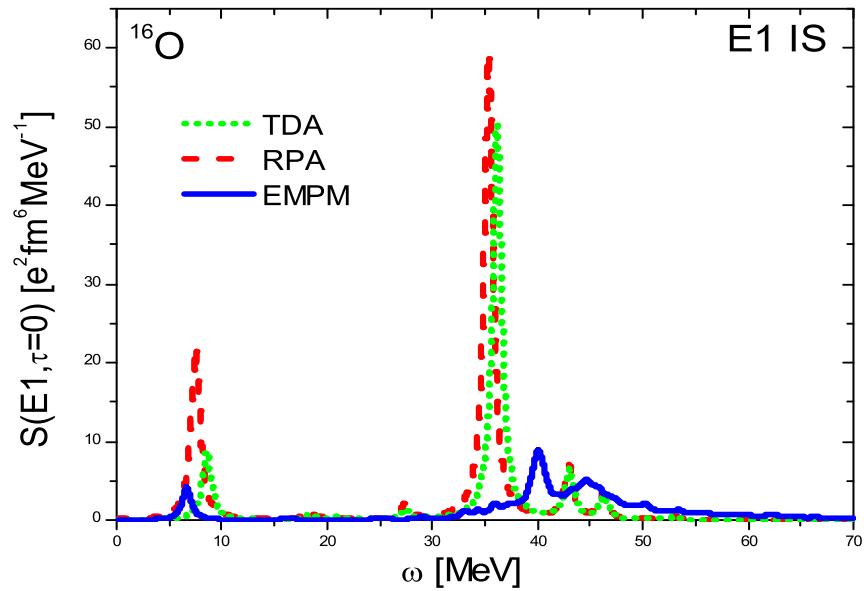
Self-consistent Green function (SCGF)

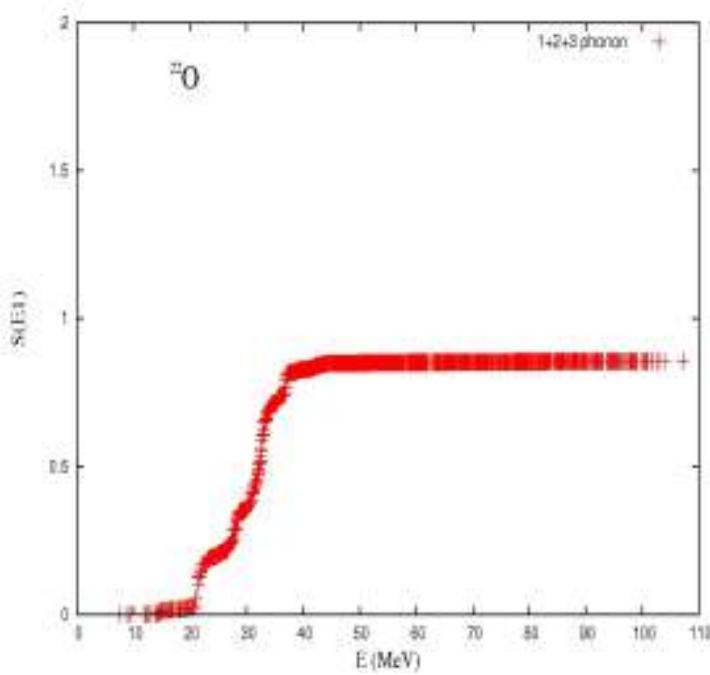
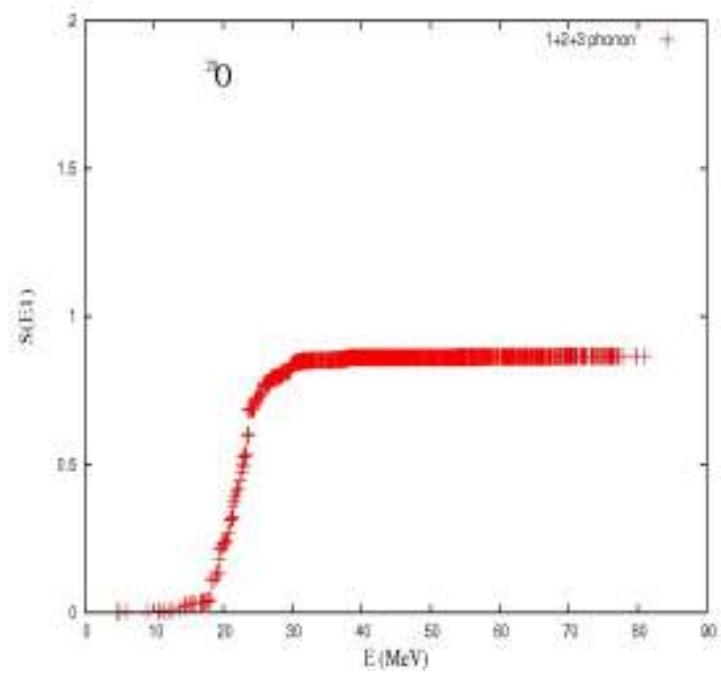
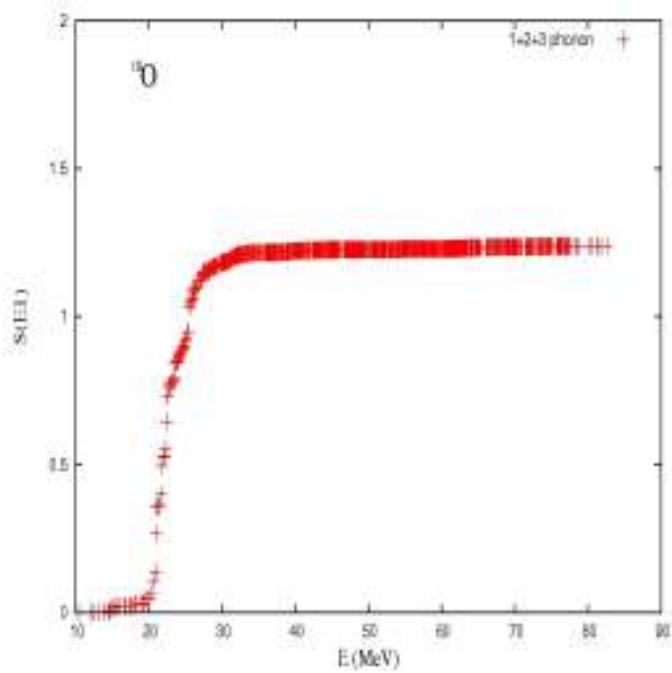
(extends RPA so as to include dressed s.p propagators and coupling to two-phonons)

C. Barbieri and W.H. Dickhoff, PRC 68, 014311 (2003);

W.H. Dickhoff and C. Barbieri, Pro. Part. Nucl. Phys. 25, 377 (2004)

# $E\lambda$ response





Running Sum